

4\_11\_08

Note Title

4/9/2008

Since  $L^2$  commutes with each component we can say

$$[L^2, \vec{L}] = 0$$

Since  $L^2$  is compatible (commutes) with each component of  $L$ , let's pick one and solve the simultaneous  $\Sigma$ -value problem:

$$L^2 f = \lambda f \quad L_z f = \mu f$$

we can only do this for observables that commute.

we'll do this by the algebraic method. i.e. ladder operators.

$$L_{\pm} = L_x \pm i L_y$$

Now can show that

$$[L_z, L_{\pm}] = \pm \hbar L_{\pm}$$

$$[L^2, L_{\pm}] = 0$$

$$\begin{aligned}
[L_z, L_{\pm}] &= [L_z, L_x \pm iL_y] \\
&= [L_z, L_x] \pm i[L_z, L_y] \\
&= i\hbar L_y \pm i(-i\hbar L_x) \\
&= \hbar(\pm L_x + iL_y) \\
&= \pm\hbar(L_x \pm iL_y) = \pm\hbar L_{\pm}
\end{aligned}$$

So lets look at

$$L_z(L_{\pm}f)$$

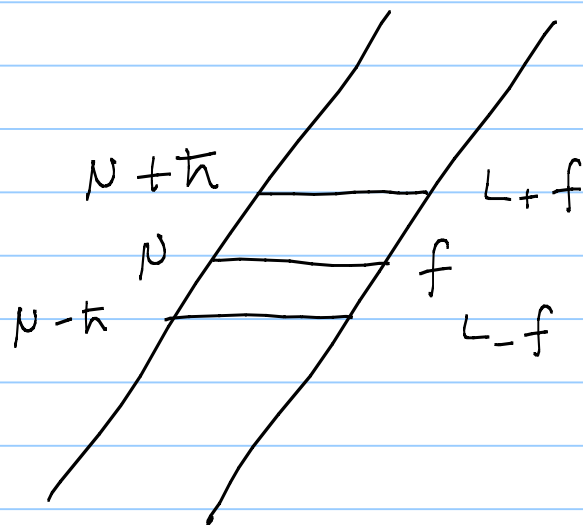
The arg. goes just like  
it did for the harm.  
oscillator

$$\begin{aligned}
L_z L_{\pm} f &= \underbrace{L_z L_{\pm} f - L_{\pm} L_z f}_{[L_z, L_{\pm}]f} + \underbrace{L_{\pm} L_z f}_{\mu f} \\
&= \pm\hbar f L_{\pm}
\end{aligned}$$

So  $L_z(L_{\pm}f) = (\mu \pm \hbar)(L_{\pm}f)$

$L_+$  is the raising operator for  
angular momentum.

$L_-$  is the lowering operator



eventually this must terminate or else we'll reach a state for which  $L_z$  exceeds  $|l|$  which cannot happen. Hence

there exists a state  $f_t$  such that

$$L_+ f_t = 0$$

So  $L_z f_t = \hbar l f_t$  for some  $l$

Now let's see what  $L^2$  does. we have assumed that

$$L^2 f_t = \lambda f_t \quad \text{now let's comp. } \lambda.$$

Somehow, if we can express  $L^2$  in terms of what we already know ( $L_z$  and  $L_{\pm}$ ) then we'll be OK.

show  $L^2 = L_+ L_- + L_z^2 + \hbar L_z$

Then  $L^2 f_t =$

$$\begin{aligned} & (L_- L_+ + L_z^2 + \hbar L_z) f_t & L_+ f_t = 0 \\ & \hbar^2 \ell^2 f_t + \hbar^2 \ell f_t \\ & = \hbar^2 \ell(\ell+1) f_t \end{aligned}$$

So

$$L^2 f_t = \hbar^2 \ell(\ell+1) f_t$$

Similarly by applying  $L_-$  we must reach a bottom state

$$L_- f_0 = 0 \quad \sim \text{for some } \bar{l}$$

$$\text{so } L_z f_0 = \hbar \bar{l} f_0$$

going through the same arg. as for  $L^2 f_l$

we can show  $L^2 f_0 = \hbar^2 \bar{l}(\bar{l}-1) f_0$

$$L^2 f_l = \hbar^2 l(l+1) f_l$$

in going up and down the ladder we're changing the  $l$ -values of  $L_z$  not  $L^2$ , so

$$\hbar^2 \bar{l}(\bar{l}-1) = \hbar^2 l(l+1) \quad \color{red}\triangleright$$

so either  $\bar{l} = l+1$  impossible for

$$\bar{l} = -l$$

Thus we conclude that the  $\Sigma$ -values of  $L_z$  are  $m\hbar$  and

$$m = -l, -l+1, \dots, l-1, l$$

integer steps.

$$2N = l \quad (\text{i.e. } l = -l + N)$$

So  $l$  itself could be integer or half integer

### Summary

The  $\Sigma$ -values of  $L_z$  are  $m\hbar$  where  $-l \leq m \leq l$  in integer steps.

but  $l$  can be integer or half int

$$L^2 f_e^m = \hbar^2 l(l+1) f_e^m$$

$$L_z f_e^m = \hbar m f_e^m$$

$$l = 0, \frac{1}{2}, 1, \dots \quad m = -l, -l+1, \dots, l$$

problem 4.18

Since  $L_{\pm} f_e^m = A_e^m f_e^{m\pm 1}$

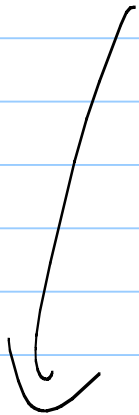
what is  $A_e^m$  if the  $f_e^m$  are to be normalized.

Lemma  $(L_{\pm})^{\dagger} = L_{\mp}$  (show this!)

recall  $L^2 = L_{\pm} L_{\mp} + L_z^2 \mp \hbar L_z$

so  $L_{\pm} L_{\mp} = L^2 - L_z^2 \pm \hbar L_z$

$\langle f_e^m | L_{\pm} L_{\mp} | f_e^m \rangle = \langle f_e^m | L^2 - L_z^2 \pm \hbar L_z | f_e^m \rangle$



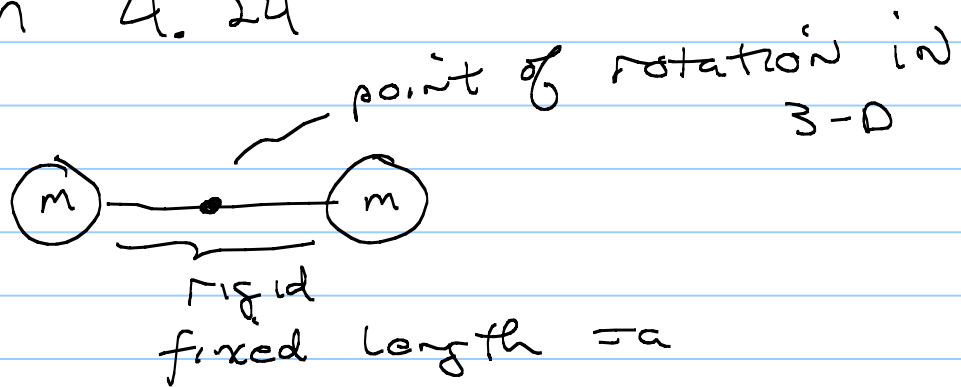
$\hbar^2 \ell(\ell+1) \langle f_e^m | f_e^m \rangle$   
 $- \hbar^2 m^2 \pm \hbar^2 m \langle f_e^m | f_e^m \rangle$

$\hbar^2 [\ell(\ell+1) - m(m\pm 1)]$

$\langle L_{\mp} f_e^m | L_{\mp} f_e^m \rangle = \langle A_e^m f_e^{m\mp 1} | A_e^m f_e^{m\mp 1} \rangle$   
 $|A_e^m|^2$

$\Rightarrow |A_e^m|^2 = \hbar^2 [\ell(\ell+1) - m(m\pm 1)]$

Problem 4.24



Total energy is 2 x kinetic  $E$  of each mass  
$$2 \left( \frac{1}{2} m v^2 \right) = m v^2$$

Angular momentum  $|\vec{L}| = 2 \frac{a}{2} m v$

$$\Rightarrow L^2 = a^2 m^2 v^2 = a^2 m H$$

$$\Rightarrow H = \frac{L^2}{a^2 m}$$

eigenvalues of  $L^2$  :  $\hbar^2 \ell(\ell+1)$

or, since we usually use  $N$  to label energies

$$E_N = \frac{1}{a^2 m} \hbar^2 N(N+1)$$

$$E_N = \frac{\hbar^2 N(N+1)}{m a^2}$$



Since  $H \propto L^2$  we know that

the stationary states are proportional to  $Y_n^m$

degeneracy of the  $n$ -th level is

$$\underline{2n + 1}$$

Problem  $L_+ Y_e^e = ?$

0

now  $L_z Y_e^e = \hbar e Y_e^e$

$L_z$  is  $-i\hbar \frac{\partial}{\partial \phi}$

so  $-i\hbar \frac{\partial}{\partial \phi} Y_e^e = \hbar e Y_e^e$

$\Rightarrow Y_e^e = f(\theta) e^{ie\phi}$