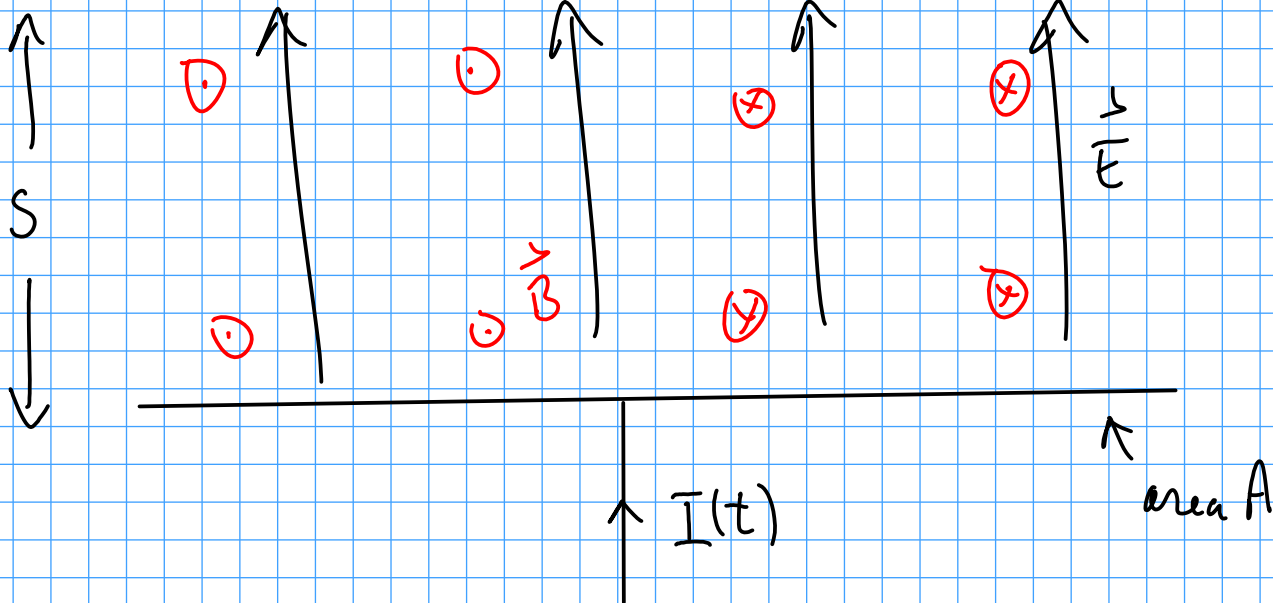
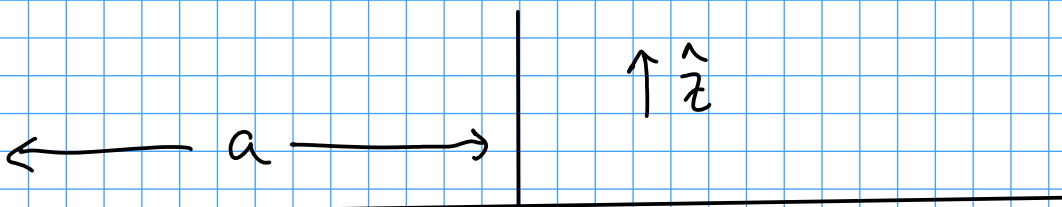
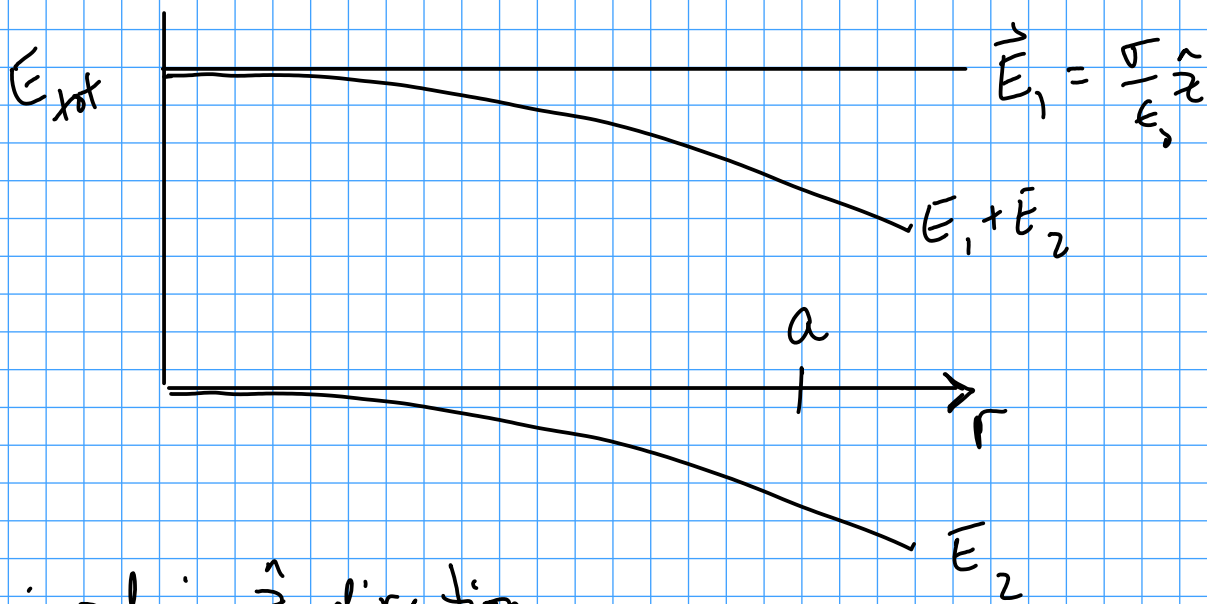
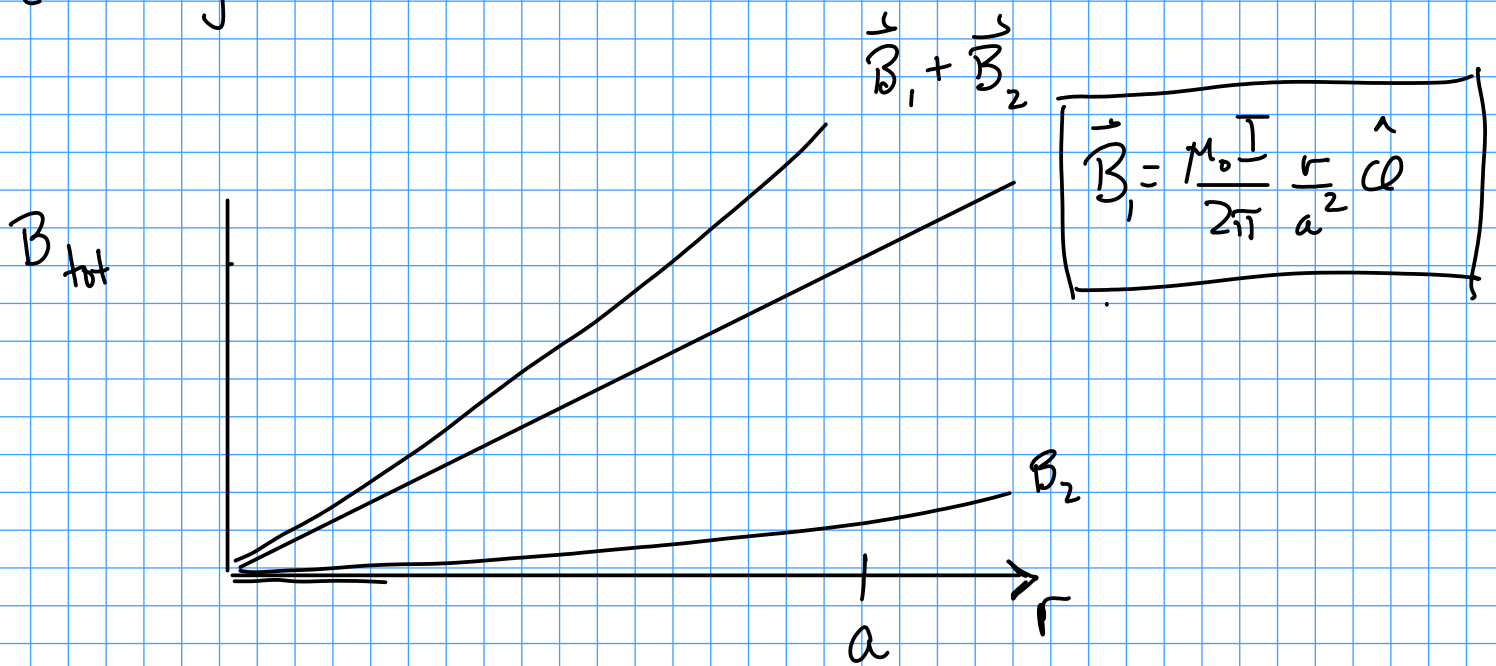


$\uparrow I(t)$





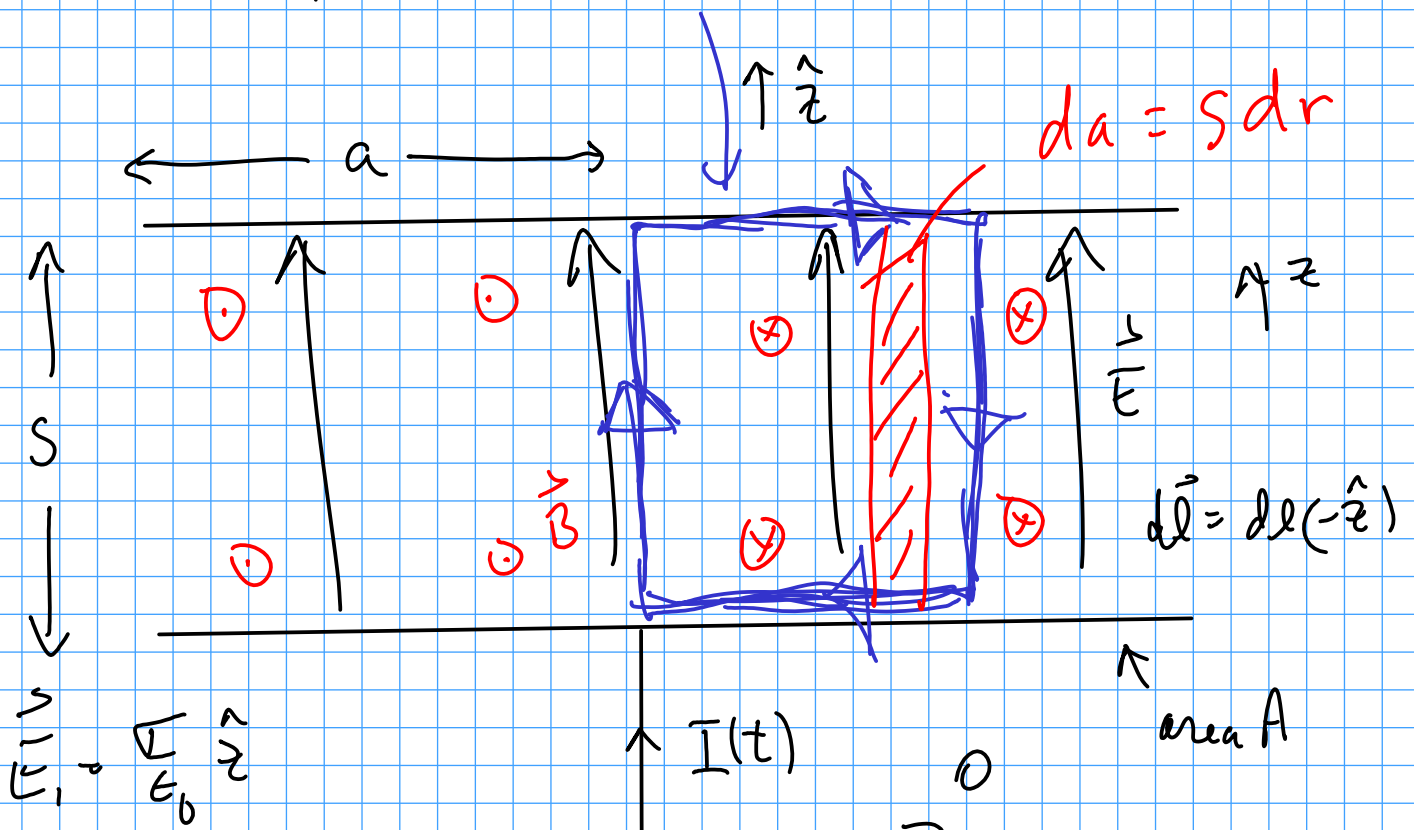
\vec{E} is only in \hat{z} direction



\vec{B} is only in $\hat{\phi}$ direction

Faraday's law

$$\oint \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \int \vec{B} \cdot d\vec{a}$$



$$\oint \vec{E} \cdot d\vec{l} = \oint (\vec{E}_1 + \vec{E}_2) \cdot d\vec{l} = \oint \vec{E}_1 \cdot d\vec{l} + \oint \vec{E}_2 \cdot d\vec{l}$$

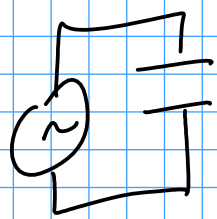
$$\int_{r=0}^a \vec{E}_2 \cdot d\vec{l} + \int \vec{E}_2 \cdot d\vec{l} = \int \vec{E}_2 \cdot d\vec{l} \text{ w/o } 180 = -E_2 s$$

$\uparrow \hat{z}$ $\uparrow -\hat{z}$

down right side

$$\vec{B}_1 = \frac{\mu_0 I}{2\pi} \frac{r}{a^2} \hat{\phi}$$

Ink Survey: Derive $\Phi_{B_1} = \int B_1 s dr$



$$\vec{E}_1(t) = E_1 e^{i\omega t}$$

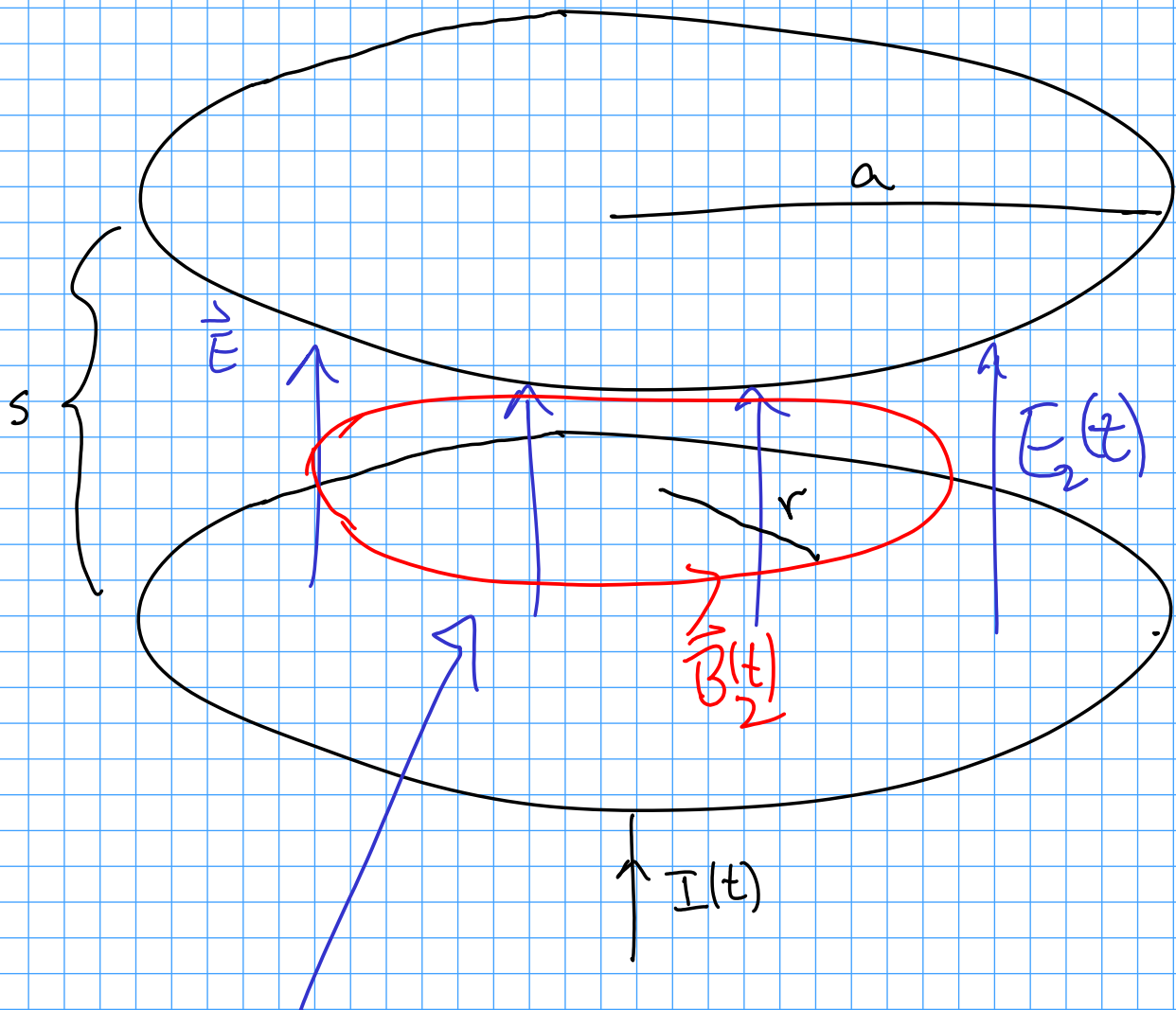
Remember: $B_1 = \frac{\mu_0 I}{2\pi a^2} r \hat{\phi} \quad \frac{1}{\epsilon_0} I = \epsilon_0 A \frac{\partial \vec{E}_1}{\partial t} = \epsilon_0 A i\omega e^{i\omega t}$

$$-\vec{E}_2 \cancel{s} = -\frac{d}{dt} \int \frac{\mu_0 r}{2\pi a^2} \epsilon_0 \overset{\pi a^2}{\cancel{a^2}} i\omega E_1 e^{i\omega t} \cancel{s} dr$$

cancel

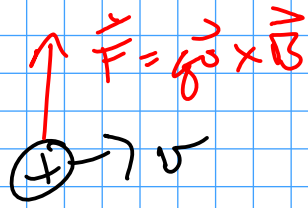
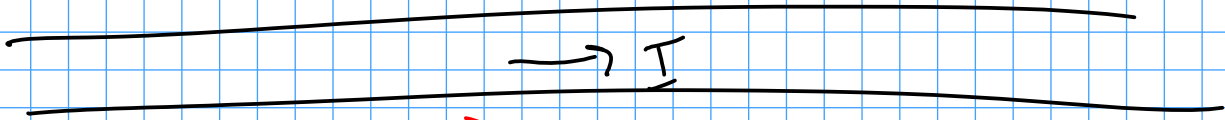
$$\vec{E}_2 = -\frac{\omega^2 r^2}{4} \mu_0 \epsilon_0 \vec{E}_1 e^{i\omega t}$$

to find \vec{B}_2 we apply



$$\oint \vec{B}_2 \cdot d\vec{l} = \mu_0 \int \vec{J}_1 \cdot d\vec{a}$$

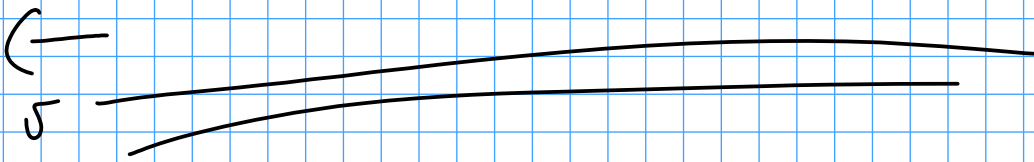
⊙ B



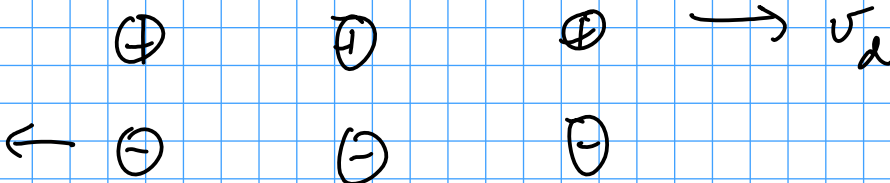
⊗ B

⊗ B

Move to the frame of particle

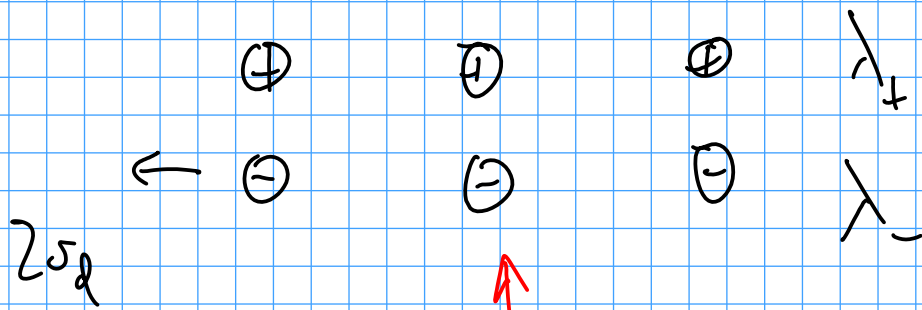


⊕ Q $v = 0$



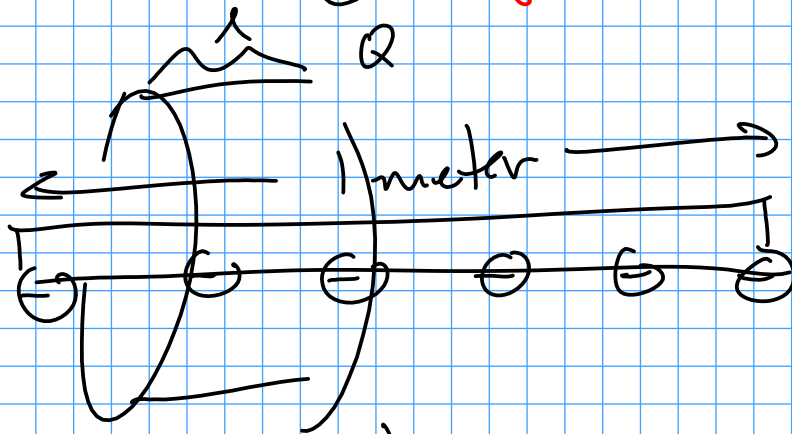
⊕ Q v_d

Frame of Q



$$F = qE$$

$|\lambda_-| > |\lambda_+|$
length contraction



$$E \cdot 2\pi r = \frac{\lambda}{\epsilon_0}$$

$$E = \frac{\lambda_+}{2\pi r \epsilon_0} + \frac{\lambda_-}{2\pi r \epsilon_0}$$