## **Resonators and stability**

Examples of resonators

Unfolded resonator and ABCD description

Stability

Ray picture

Gaussian beam picture

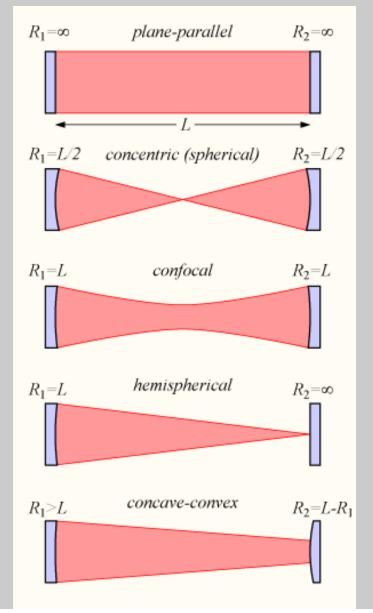
2 mirror resonators and the stability map

Analysis of resonators

beam sizes

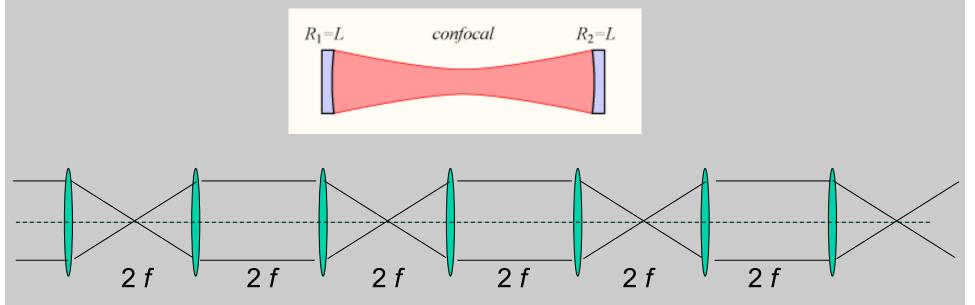
### Resonators

- Resonators provide feedback for the photons to build up by passing through the gain medium
- Curved mirrors are typically used to control the beam size inside the gain medium
- Types of resonators
  - Many resonators have more than two mirrors, but most can be mapped onto a two-mirror system.



### **Periodic lens model**

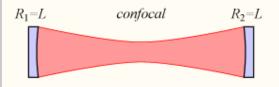
 A resonator can be "unfolded" by modeling the curved mirrors as ideal lenses



- Are there rays that will stay confined?
- If so, resonator is *stable*.

## **Resonator ABCD model**

- Build a ABCD matrix model of the periodic lens sequence
  R<sub>1</sub>=L confocal
  - First get a round-trip matrix



$$M_{RT}(f,L) = \begin{pmatrix} A & B \\ C & D \end{pmatrix} = M_L(f) \cdot M_T(L) \cdot M_L(f) \cdot M_T(L)$$
$$= \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

- Free to choose starting point
- Focal length **f** and mirror separation **L** can vary

$$\begin{pmatrix} r_2 \\ r'_2 \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} A & B \\ C & D \end{pmatrix} \begin{pmatrix} r_0 \\ r'_0 \end{pmatrix}$$

After 2 round trips

#### **Resonator stability: ray picture**

- Will a ray stay trapped?
- Look at whether  $r_n$  and  $r'_n$  stay finite as n goes to infinity  $\begin{pmatrix} r_n \\ r'_n \end{pmatrix} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}^n \begin{pmatrix} r_0 \\ r'_0 \end{pmatrix}$
- Method:  $M_{RT} = \begin{pmatrix} A & B \\ C & D \end{pmatrix} = U \begin{pmatrix} \lambda_a & 0 \\ 0 & \lambda_b \end{pmatrix} U^{-1}$ 
  - Diagonalize matrix:

$$- \text{ then } M_{RT}^{n} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}^{n} = U \begin{pmatrix} \lambda_{a} & 0 \\ 0 & \lambda_{b} \end{pmatrix} U^{-1} U \begin{pmatrix} \lambda_{a} & 0 \\ 0 & \lambda_{b} \end{pmatrix} U^{-1} \cdots$$
$$= U \begin{pmatrix} \lambda_{a} & 0 \\ 0 & \lambda_{b} \end{pmatrix}^{n} U^{-1} = U \begin{pmatrix} \lambda_{a}^{n} & 0 \\ 0 & \lambda_{b}^{n} \end{pmatrix} U^{-1}$$

# **Stability condition**

• The ray will stay trapped (stable resonator) if

$$|\lambda_a| \le 1$$
  $|\lambda_b| \le 1$   $\frac{1}{|\lambda_a|} \le 1$   $\frac{1}{|\lambda_b|} \le 1$  Reverse propagation

- Therefore matrix eigenvalues must satisfy  $|\lambda_a| = |\lambda_b| = 1$
- Property of ABCD: det  $M_{RT} = \lambda_a \lambda_b = 1$

 $\therefore \lambda_a = e^{i\theta}, \lambda_b = \lambda_a^{*}$ 

• Trace of M is invariant upon rotation of matrix:

$$\operatorname{Tr} M_{RT} = \lambda_a + \lambda_b = A + D$$
$$= e^{i\theta} + e^{-i\theta} = 2\cos\theta$$

• Finally stability condition is:

$$-1 \le \frac{A+D}{2} \le 1$$

## Some properties of ABCD matrices

- 1. Determinant = 1 if start and end points are in the same medium (same refr. Index)
  - Example:  $\begin{pmatrix} 1 & 0 \\ -\frac{1}{f} & 1 \end{pmatrix} \begin{pmatrix} 1 & L \\ 0 & 1 \end{pmatrix}$ - Counter example: dielectric interface  $\begin{pmatrix} 1 & 0 \\ 0 & n_1/n_2 \end{pmatrix}$
  - Therefore, det M = 1, but note that eigenvalues can be real <u>or</u> complex
- 2. Complex eigenvalues are of the form  $e^{\pm i\theta}$ 
  - Outside of stability range, eigenvalues are <u>real</u>

 $\lambda_a = 1 / \lambda_b$  Tr  $M_{RT} = \lambda_a + 1 / \lambda_a > 2$  if  $\lambda_a > 1$ 

3. M is not necessarily unitary (where  $M^{-1} = M^{\dagger}$ )

# Stability for *Gaussian* beams in resonators

- A stable resonator mode is one that repeats itself on each round trip
  - Amplitude and phase are matched  $\therefore q_{n+1} = q_n$

$$q_1 = \frac{Aq_0 + B}{Cq_0 + D} = q_0 \rightarrow Aq_0 + B = q_0(Cq_0 + D)$$

$$\rightarrow 0 = Cq_0^2 + (D - A)q_0 - B$$

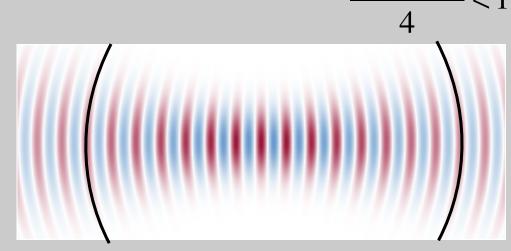
$$q_0 = \frac{(A-D)}{2C} \pm \frac{1}{2C} \sqrt{(A-D)^2 + 4BC}$$

- Since  $\frac{1}{q_0} = \frac{1}{R} - i \frac{\lambda}{\pi w^2}$   $q_0$  must be complex (w is finite)

$$\therefore \left(A-D\right)^2 + 4BC < 0$$

# Stability for Gaussian beams in resonators

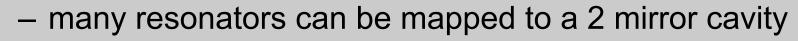
- We know:  $(A-D)^2 + 4BC < 0$
- And, since det(M) = 1 AD BC = 1  $(A - D)^2 + 4BC = (A - D)^2 + 4(AD - 1)$   $= A^2 - 2AD + D^2 + 4AD - 4$  $= (A + D)^2 - 4 < 0$
- Stability condition:  $(A+D)^2 < 1$

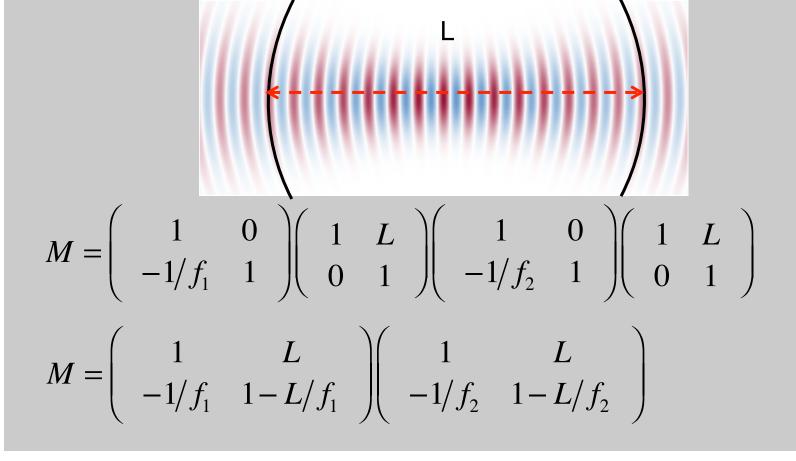


If this condition is satisfied, curvature of each end mirror matches wavefront curvature.

### 2 mirror cavity stability

• Important example





### Stability for 2 mirror resonator

• Stability condition:  $\frac{(A+D)^2}{4} < 1 \rightarrow -1 < \frac{A+D}{2} < 1$ 

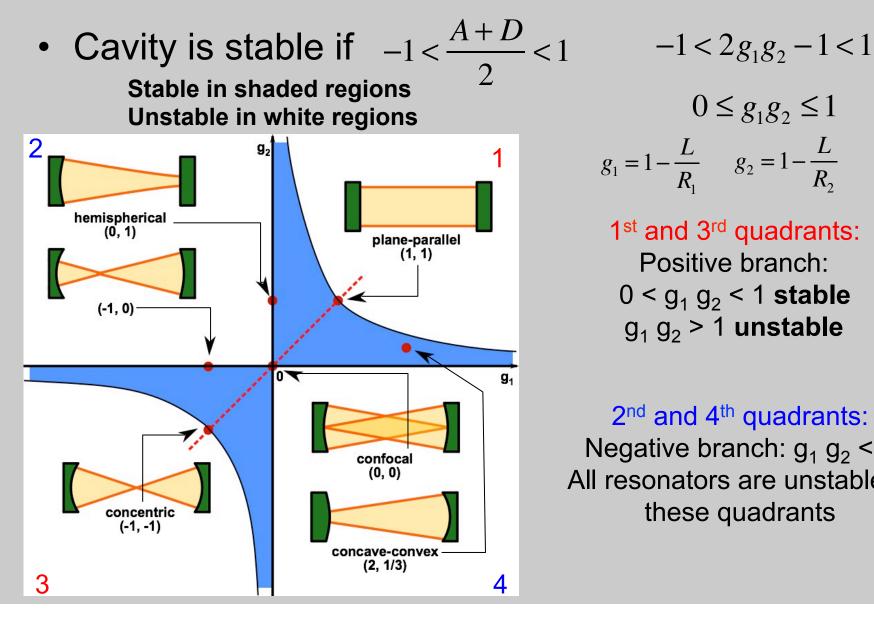
- Evaluate A and D from round-trip matrix

$$M = \begin{pmatrix} 1 & L \\ -1/f_1 & 1 - L/f_1 \end{pmatrix} \begin{pmatrix} 1 & L \\ -1/f_2 & 1 - L/f_2 \end{pmatrix}$$

 $A = 1 - L/f_2$   $D = -L/f_1 + (1 - L/f_1)(1 - L/f_2)$   $f_1 = R_1/2$  $f_2 = R_2/2$ 

$$\frac{A+D}{2} = \frac{1}{2} \left( 1 - \frac{2L}{R_2} - \frac{2L}{R_1} + 1 - \frac{2L}{R_1} - \frac{2L}{R_2} + \frac{4L^2}{R_1R_2} \right)$$
$$= 1 - \frac{2L}{R_1} - \frac{2L}{R_2} + \frac{2L^2}{R_1R_2} = 2 \left( 1 - \frac{L}{R_1} \right) \left( 1 - \frac{L}{R_2} \right) - 1 \equiv 2g_1g_2 - 1$$

#### 2 mirror stability and the stability map



 $0 \le g_1 g_2 \le 1$  $g_1 = 1 - \frac{L}{R_1}$   $g_2 = 1 - \frac{L}{R_2}$ 

1<sup>st</sup> and 3<sup>rd</sup> quadrants: Positive branch:  $0 < g_1 g_2 < 1$  stable  $g_1 g_2 > 1$  unstable

2<sup>nd</sup> and 4<sup>th</sup> quadrants: Negative branch:  $g_1 g_2 < 0$ All resonators are unstable in these quadrants

# **Boundaries of stability** $g_1 = 1 - \frac{L}{R_1}$ $g_2 = 1 - \frac{L}{R_2}$

• Easily identified stable resonators are actually at edge of stability

