

Reading: Today: 12.2  
Tomorrow: 12.3

Last time: covariant / contravariant  
vectors,

$$x^{\mu} = (ct, x, y, z)$$

↑    ↑    ↑    ↑  
 $x^0$   $x^1$   $x^2$   $x^3$

Contravariant  $\bar{x}^{\nu} = \frac{\partial x^{\nu}}{\partial x^{\mu}} x^{\mu}$

any old contravariant vector

$$a^{\nu} = \frac{\partial x^{\nu}}{\partial x^{\mu}} a^{\mu}$$

for covariant vectors

$$\bar{a}_{\nu} = \frac{\partial x^{\mu}}{\partial x^{\nu}} a_{\mu}$$

Use Lorentz transform for  $\bar{S}$  going rel.  $v$  w respect to  $S$  in  $x$ -direction

$$\begin{aligned} \bar{x} &= \gamma(x-vt) & \text{Find } \frac{\partial \bar{x}^\nu}{\partial x^\mu} \text{ matrix} \\ \bar{y} &= y \\ \bar{z} &= z \\ \bar{t} &= \gamma(t - \frac{v}{c}x) \end{aligned}$$

$$\text{In this case: } \frac{\partial \bar{x}^\nu}{\partial x^\mu} = \begin{pmatrix} \gamma & -\gamma v/c & 0 & 0 \\ -\gamma v/c & \gamma & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

$$\bar{x}^\nu = \frac{\partial \bar{x}^\nu}{\partial x^\mu} x^\mu$$

↑  
If this guy were  $\Lambda_{\mu\nu}$   
→  $\Lambda_{\mu\nu} x^\mu = \bar{x}_\nu$

So what about the operator  $\frac{\partial}{\partial x^\mu}$ ?

$$\partial_{\bar{x}^\nu} = \bar{x}_i^\nu - \bar{x}_0^\nu \rightarrow \frac{\partial}{\partial x^\mu} \text{ is covariant}$$

A lot of times you'll see people write

$$\frac{\partial}{\partial x^\mu} = \partial_\mu = \frac{\partial}{\partial x^0} \hat{x}^0 + \frac{\partial}{\partial x^1} \hat{x}^1 + \frac{\partial}{\partial x^2} \hat{x}^2 + \frac{\partial}{\partial x^3} \hat{x}^3$$

↑  
This operator replaces  $\vec{\nabla}$  in 30.

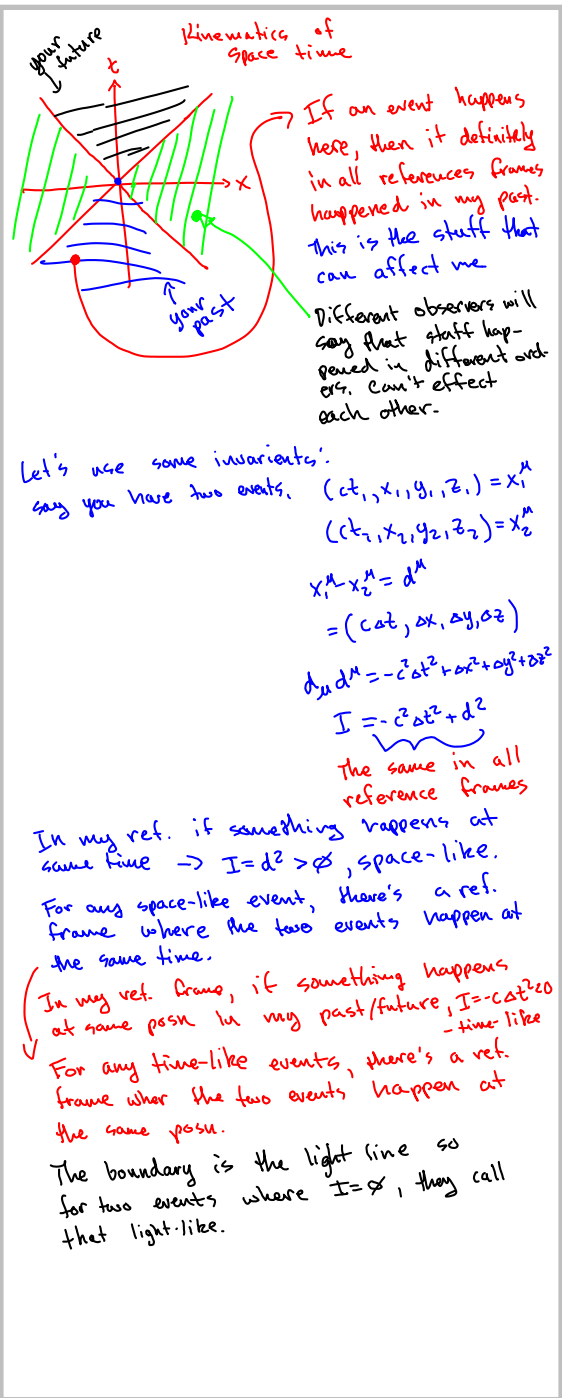
$$\text{Find } \partial_\mu \partial^\mu : \partial_\mu = \left( \frac{1}{c} \frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

$$\partial^\mu = g_{\mu\nu} \partial^\nu = \left( -\frac{1}{c} \frac{\partial}{\partial t}, \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right)$$

$$\partial_\mu \partial^\mu = g_{\mu\nu} \partial^\nu \partial^\mu = \underbrace{-\frac{1}{c^2} \frac{\partial^2}{\partial t^2} + \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}}_{\partial_\mu \partial^\mu}$$

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) V = \frac{-\rho}{\epsilon_0}$$

$$\left( \nabla^2 - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} \right) \vec{A} = -\mu_0 \vec{J}$$



If an event happens here, then it definitely in all reference frames happened in my past. this is the stuff that can affect me

Different observers will say that stuff happened in different order. Can't affect each other.

Let's use some invariants:

say you have two events,  $(ct_1, x_1, y_1, z_1) = x_1^M$

$(ct_2, x_2, y_2, z_2) = x_2^M$

$$x_1^M - x_2^M = d^M$$

$$= (c\Delta t, \Delta x, \Delta y, \Delta z)$$

$$d_{\mu\nu} d^{\mu\nu} = -c^2 \Delta t^2 + \Delta x^2 + \Delta y^2 + \Delta z^2$$

$$I = -c^2 \Delta t^2 + d^2$$

The same in all reference frames

In my ref. if something happens at same time  $\rightarrow I = d^2 > 0$ , space-like.

For any space-like event, there's a ref. frame where the two events happen at the same time.

In my ref. frame, if something happens at same posn in my past/future,  $I = -c\Delta t^2 < 0$  -time-like

For any time-like events, there's a ref. frame where the two events happen at the same posn.

The boundary is the light line so for two events where  $I = 0$ , they call that light-like.

Proper time:  $d\tau = \sqrt{1 - \frac{u^2}{c^2}} dt = \frac{1}{\gamma} dt$

Let's say you have a moving particle in your lab moving w/ speed  $u$ .

change in time in part. frame

change for lab frame

Proper velocity:  $\vec{v} = \gamma \vec{u}$

$\frac{d\vec{b}}{dt}$  ← lab frame

$\vec{u} = \frac{d\vec{b}}{dt}$

$\frac{d\vec{b}}{d\tau}$  ← object frame

No 4-vector exists for  $\vec{u}$ . But there is one for  $\vec{v}$

$$\eta^\mu = \begin{pmatrix} \gamma c \\ \gamma \vec{u} \end{pmatrix} \quad \eta^\mu = \frac{\partial \bar{x}^\mu}{\partial x^\nu} \eta^\nu$$

That also means that

$$\eta_\mu \eta^\mu \text{ is invariant}$$

$$-\gamma^2 c^2 + \gamma^2 u^2 = \text{invariant.}$$

If I multiply by  $m^2 c^2$

$$-\gamma^2 m^2 c^4 + \gamma^2 m^2 u^2 c^2$$

Next time we'll see that

$$\vec{p} = \gamma m \vec{v} ; E = \gamma m c^2$$

↑ momentum      ↑ energy

$$-E^2 + (pc)^2$$