Day 13: Separation of variables in cylindrical coords

Laplace's eqn in cylindrical coords looks like: Ask them to separate the variables?

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial V}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 V}{\partial \phi^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

As usual, we'll cut ourselves a break and only consider "long" objects, so that there's no z-dependence in V and we can cut down to

$$\frac{1}{r}\frac{\partial}{\partial r}\left(r\frac{\partial V}{\partial r}\right) + \frac{1}{r^2}\frac{\partial^2 V}{\partial \phi^2} = 0$$

Also as usual, hope for a solution of the form $V(r,\phi) = R(r)\,\Phi(\phi)$

Plug & crank to get

$$\Phi\frac{1}{r}\frac{d}{dr}\left(r\frac{d(R)}{dr}\right) + \frac{1}{r^2}R\frac{d^2\Phi}{d\phi^2} = 0$$

$$\Phi\left[\frac{1}{r}\frac{dR}{dr} + \frac{d^2R}{dr^2}\right] + \frac{1}{r^2}R\frac{d^2\Phi}{d\phi^2} = 0$$ Multiply by $\frac{r^2}{\Phi R}$ to separate

$$\frac{1}{R}\left[r\frac{dR}{dr} + r^2\frac{d^2R}{dr^2}\right] + \frac{1}{\Phi}\frac{d^2\Phi}{d\phi^2} = 0$$

So $\frac{1}{\Phi}\frac{d^2\Phi}{d\phi^2} = k_1^2$, $\frac{1}{R}\left[r\frac{dR}{dr} + r^2\frac{d^2R}{dr^2}\right] = k_2^2$ with $k_1^2 + k_2^2 = 0$

The Φ solutions need to be periodic (ϕ coordinate goes from 0 to 2π), so they're the sines/cosines: $(k_1^2 \leq 0)$

$$\Phi(\phi) = C_n\cos(n\phi) + D_n\sin(n\phi)$$

Note the periodicity also forces k_1 to be integer, hence the n $\cos(0) = \cos(2\pi k_1)$

That leaves (since $k_2^2 = n^2 \geq 0$)

$$r\frac{dR}{dr} + r^2\frac{d^2R}{dr^2} - n^2R = 0$$ This is in our bag of semi-standard Diff Eqs, and the general solution is

$$R(r) = A_n r^n + B_n r^{-n}$$ for any positive integer n (easy-ish guess & check)

When $n=0$ we get a special case:

$$r\frac{d}{dr}\left(r\frac{dR}{dr}\right) = 0$$ With solution $A\ln r + B$

(write as total derivative $r\frac{d}{dr}\left(r\frac{dR}{dr}\right) = 0$)

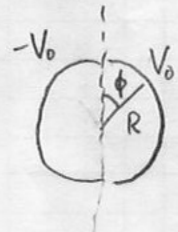
The full meal deal is

$$V(r,\phi) = A\ln(r) + B + \sum_{n=1}^{\infty}\left(A_n r^n + B_n r^{-n}\right)\underbrace{\left(C_n \cos n\phi + D \sin n\phi\right)}$$

Zonal harmonics, or, as normal people would say, sines and cosines

Note that, as usual, you can often exclude big chunks of this based on symmetry or other considerations.

Example: Double hemi-cylinder

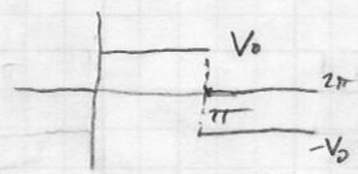


Conducting cylinder of radius R split into two halves held at $-V_0$ and $+V_0$. Lets hack away at V (everywhere)

Ditch B. Poor B.

Ditch A. This object is net neutral, so as $r \to \infty$, we don't expect $V \to \infty$ (or at $r=0$, for that matter)

Now, the potential BC <u>as I've drawn</u> <u>it</u> is odd in ϕ, so we need ϕ-functions that respect that. Lose the cosines. (imagine making it periodic)



Here's a twist: r^n blows up as $r \to \infty$ (which we don't think we want) and r^{-n} blows up as $r \to 0$ (which we're <u>really</u> <u>quite</u> <u>sure</u> we don't want)

But we need some r-dependence, so we define V piecewise:

One V for $r \geq R$: $\quad V(r,\phi)_{outside} = \sum_{n=1}^{\infty} B_n r^{-n} \sin(n\phi)$

Another V for $r \leq R$: $\quad V(r,\phi)_{inside} = \sum_{n=1}^{\infty} A_n r^n \sin(n\phi)$

We'll need to make sure this is continuous across $r=R$. For that to happen we need

$$B_n R^{-n} = A_n R^n$$ For all n (a rather harsh restriction)

Each of these is constant and equal to the other, so we can use a little shortcut gimmick and say

$$B_n R^{-n} = A_n R^n = C_n \quad \Rightarrow \quad B_n = C_n R^n$$

$$A_n = C_n R^{-n}$$

Leaving us with only one set of constants to fix:

$$V(r,\phi)_{out} = \sum_{n=1}^{\infty} c_n \left(\frac{R}{r}\right)^n \sin(n\phi)$$

$$V(r,\phi)_{in} = \sum_{n=1}^{\infty} c_n \left(\frac{r}{R}\right)^n \sin(n\phi)$$

Now we use Fourier's trick at $r=R$, where $V(R,\phi) = \sum_{n=1}^{\infty} c_n \sin(n\phi)$

$$\int_0^{2\pi} V(R,\phi)\sin(m\phi)\,d\phi = \sum_{n=1}^{\infty} c_n \int_0^{2\pi} \underbrace{\sin(n\phi)\sin(m\phi)\,d\phi}_{\pi\delta_{nm}} = c_m \cdot \pi$$

$V(R,\phi)$ is V_0 from 0 to π, $-V_0$ from π to 2π

$$\int_0^{\pi} V_0 \sin(m\phi) + \int_{\pi}^{2\pi} -V_0 \sin(m\phi) = V_0 \cdot \left[-\frac{1}{m}\cos(m\phi)\Big|_0^{\pi} + \frac{1}{m}\cos(m\phi)\Big|_{\pi}^{2\pi} \right]$$

$$= V_0\left[-\frac{1}{m}\cos(m\pi) + \frac{1}{m}\cos(0) + \frac{1}{m}\cos(2\pi m) - \frac{1}{m}\cos(m\pi)\right]$$

$$= V_0\left[-\frac{2}{m}\cos(m\pi) + \frac{2}{m}\right]$$

$$= \frac{2V_0}{m}\left[-\cos(m\pi) + 1\right]$$

For $m = 0, 2, 4,$ etc $\cos m\pi = 1$ and this zeroes out

For $m = 1, 3,$ etc $\cos(m\pi) = -1$ and we get $\frac{4V_0}{(2j+1)}$

$$\Rightarrow c_j = \frac{4V_0}{\pi(2j+1)} \quad \text{for all } j$$

Putting it together, $V(r,\phi)_{out} = \frac{4V_0}{\pi} \sum_{j=1}^{\infty} \frac{1}{(2j+1)} \left(\frac{R}{r}\right)^{(2j+1)} \sin((2j+1)\phi)$

$$V(r,\phi)_{in} = \frac{4V_0}{\pi} \sum_{j=1}^{\infty} \left(\frac{1}{2j+1}\right) \left(\frac{r}{R}\right)^{(2j+1)} \sin((2j+1)\phi)$$

Note that I set up my coordinates slightly different than the book did, and so get a slightly different answer in ϕ. It's worth pondering how these can be equivalent, since nothing changed physically.