

$$t_r = t - \frac{r}{c} \quad \text{source \& rec. are fixed}$$

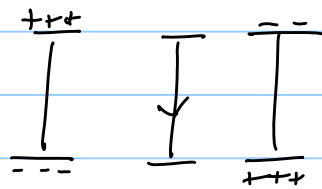
Note Title

7/3/2006

Ex:

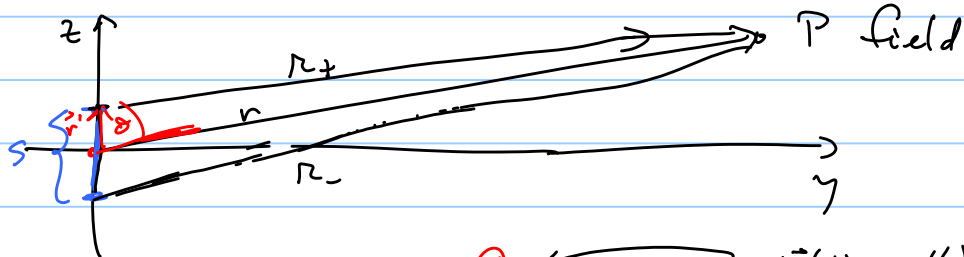


antenna



$$q(t) = q_0 \cos \omega t$$

$$I = \frac{dq}{dt}$$



$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\vec{r}', t_r)}{r} d\tau'$$

$$\rho(\vec{r}', t) = q(t) \delta(\vec{r}' - \frac{s}{2} \hat{z}) - q(t) \delta(\vec{r}' + \frac{s}{2} \hat{z})$$

$$V(\vec{r}, t) = \frac{1}{4\pi\epsilon_0} \left[q_0 \frac{\cos[\omega(t - \frac{r_+}{c})]}{r_+} - q_0 \frac{\cos[\omega(t - \frac{r_-}{c})]}{r_-} \right]$$

Approx: far away from dipole $s \ll r$
 $\lambda \ll r$ far field approx

$$r_+^2 = \vec{r}_+ \cdot \vec{r}_+ = r^2 + r'^2 - 2\vec{r} \cdot \vec{r}' \quad \vec{r}' = \frac{s}{c} \hat{j}$$

$$|\vec{r}_+| \approx r \left(1 - \frac{s}{2r} \cos \theta \right) \quad \text{: keep only terms 1st order in } \frac{s}{r}$$

$$V(r, \theta, t) = - \frac{P_0 \omega}{4\pi \epsilon_0 c} \frac{\cos \theta}{r} \sin \omega \left(t - \frac{r}{c} \right)$$

$\omega \left(t - \frac{r}{c} \right) \left(1 - \frac{s}{2r} \cos \theta \right) \omega$
 $v = c$
 $\frac{v}{c} = \frac{s}{r}$

P_0 is dipole moment = $q_0 s$

Vector potential $\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \int \frac{I(\vec{r}', t_r)}{r} d\vec{r}'$

$$= \frac{\mu_0}{4\pi} \int \frac{-q_0 \omega \sin[\omega(t - \frac{r}{c})] \hat{z}}{r} dz' \quad \int \delta(x) dx = 1$$

To 1st order in $\frac{s}{r}$

$$\vec{A}(\vec{r}, t) = \frac{\mu_0}{4\pi} \frac{-q_0 \omega \sin \omega \left(t - \frac{r}{c} \right)}{r} \int_{-s/2}^{s/2} dz' \hat{z}$$

In far field this appears to be a spherical wave

$$\hat{k} = \cos \theta \hat{r} - \sin \theta \hat{\theta} \quad \text{find } \vec{E} \text{ \& } \vec{B} = \nabla \times \vec{A}$$

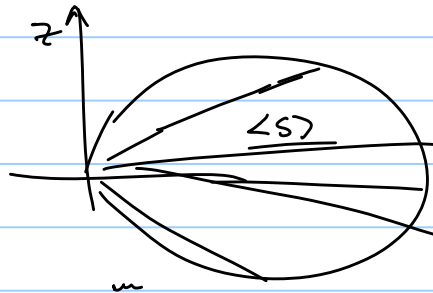
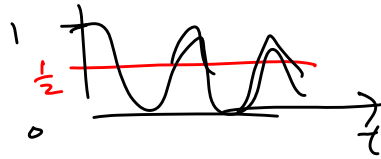
$-\nabla \cdot \frac{\partial \vec{A}}{\partial t}$

Find $\langle \frac{\vec{E} \times \vec{B}}{\mu_0} \rangle$

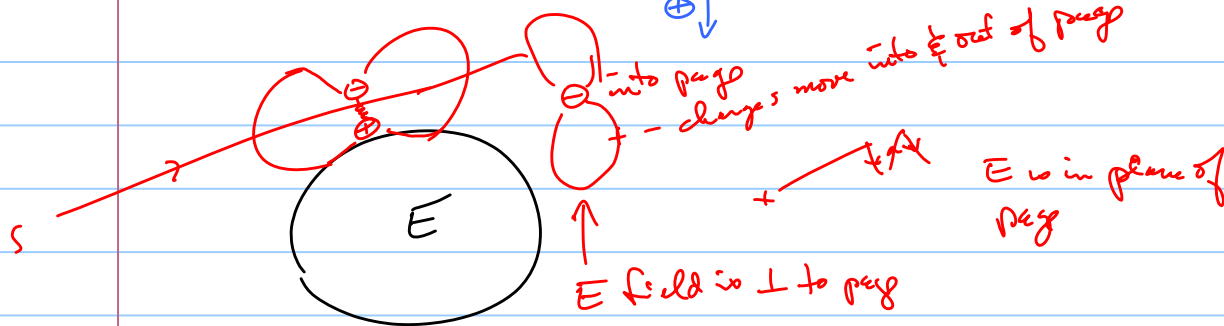
$$\langle S \rangle = \frac{\mu_0 p_0^2 \omega^4}{32 \pi^2 c}$$

$$\cos^2(\omega t \dots)$$

$$\frac{\sin^2 \theta}{r^2}$$



Lookout into \rightarrow atoms



Larmor formula for power radiated

$$P = \int \vec{S} \cdot d\vec{a} = \frac{\mu_0}{6\pi c} \ddot{p}^2$$

p is dipole moment

$\vec{P} = q \vec{s}$ for two charges $\vec{P} = \int_{\text{all } \vec{r}'} dq \vec{r}' = \int \rho d\vec{r}' \vec{r}'$

1-D $\vec{P} = \int \lambda d\vec{r}'$

$\lambda = \lambda_0 \sin \phi$

$d\vec{r}' = b d\phi$

$\vec{r}' = b \hat{r}$

0

$\int_0^{2\pi} \lambda_0 \sin \phi b d\phi b \hat{r}$

cannot come outside integral because it depends on ϕ

2π

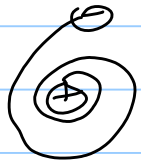
$\int_0^{2\pi} \lambda_0 \sin \phi b d\phi [b \cos \phi \hat{x} + b \sin \phi \hat{y}] = \vec{P}$

$\phi(t)$



$P_{\text{rad}} = \frac{a_1}{r^4}$

$$U_{\text{atom electric}} = U_{\text{kin}} + U_{\text{potential energy}} = -\frac{a_2}{r(t)}$$



$v(t)$

$$-\frac{dU_{\text{tot}}}{dt} = -\frac{dU}{dr} \frac{dr}{dt} = P_{\text{radiation}} = \text{Larmor for}$$

loss of energy

\therefore

$a_{\text{rad}} = \frac{v^2}{r}$ circular motion

$$F = ma = m \frac{v^2}{r}$$

Coulomb

