

These are the dynamic Maxwell equations

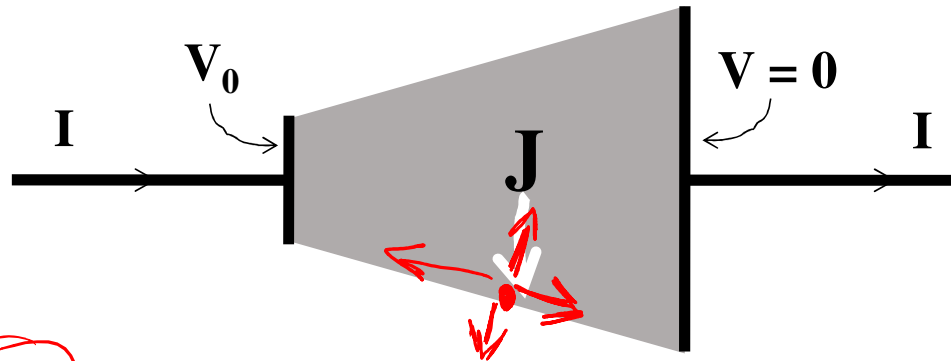
$$\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0} \quad \nabla \cdot \vec{B} = 0$$

$$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \quad \nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

What can I do to make them be the dynamic Maxwell equations *in matter*?

- A) Leave them alone, but attend carefully to what  $\rho$  and  $J$  are. 6
- B) Replace the  $\epsilon_0$  and  $\mu_0$  with  $\epsilon$  and  $\mu$ , and let the sources be free sources. 0
- C) Make a variety of substitutions involving  $\epsilon$ ,  $\mu$ ,  $H$ , and  $D$ . 2
- D) A and C 27
- E) A, B, and C 7

Inside this resistor setup, what can you conclude about the current density  $\mathbf{J}$  near the side walls?



- A) Must be exactly parallel to the wall 18
- B) Must be exactly perpendicular to the wall 3
- C) Could have a mix of parallel and perpendicular components 21
- D) No obvious way to decide

Let's get some more perspective on the previous. Starting from the continuity equation for current,  $\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$ , derive a general boundary condition on the perpendicular component of  $J$ .

A)  $J_{\text{perp},1} - J_{\text{perp},2} = 0$

B)  $J_{\text{perp},1} - J_{\text{perp},2} = -\frac{\partial \rho}{\partial t}$

C)  $J_{\text{perp},1} - J_{\text{perp},2} = -\frac{\partial \sigma}{\partial t}$

D)  $J_{\text{perp},1} - J_{\text{perp},2} = -\rho$

E) Something else

Let's make use of that here.

