

Radiation w/ velocity // acceleration

$$\vec{E}_a = \frac{e}{c^2 K^3 R^2} \vec{R} ((\vec{R} - \vec{\beta} R) \times \vec{a})$$

\downarrow
 0 since here $\vec{\beta} \times \vec{a} = 0$

now $\vec{S}_a = \vec{S}_a(\beta=0) \cdot K^{-6}$

but this is power radiated from charge:

$$\frac{dE_{\text{tot}}}{dt} = - \int \vec{S}_a \cdot d\vec{A}$$

we want radiated power measured at distant point,

$$P = - \frac{dE_{\text{tot}}}{dt'} = - \frac{dE}{dt} \frac{dt}{dt'} \quad t' = t - R(t)/c$$

↳ retarded time

$$\frac{dt}{dt'} = 1 + \frac{1}{c} \frac{d}{dt'} |\vec{r} - \vec{r}_e(t)| = 1 + \sum \frac{u_i}{c} \frac{d}{dx_{ei}} |\vec{r} - \vec{r}_e|$$

$$|\vec{r} - \vec{r}_e| = \sqrt{r^2 + r_e^2 - 2\vec{r} \cdot \vec{r}_e}$$

$$\frac{d}{dx_{ei}} |\vec{r} - \vec{r}_e| = \frac{1}{2} \frac{1}{\sqrt{\dots}} (2x_{ei} - 2x_{ei})$$

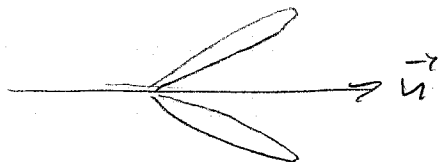
$$\sum \beta_i \frac{d}{dx_{ei}} |\vec{r} - \vec{r}_e| = \sum \frac{\beta_i (x_{ei} - x_i)}{R} = - \frac{\vec{\beta} \cdot \vec{R}}{R}$$

$$\frac{dt}{dt'} = 1 - \frac{\vec{\beta} \cdot \vec{R}}{R} = K = 1 - \beta \cos \theta$$

↳ measured rel to \vec{u}

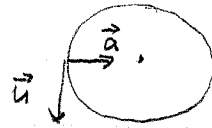
now we have K^{-5} in $dP/d\Omega$

this directs lobes toward \vec{u}

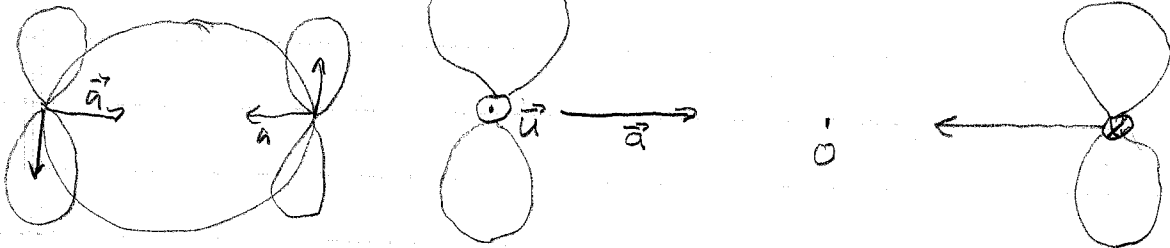


Synchrotron radiation - circular orbit.

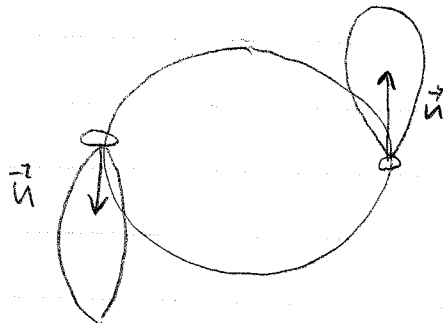
$$\vec{u} \perp \vec{a}$$



At low speeds, Larmor formula w/ θ relative to \vec{a}



higher speeds: push lobes forward, toward \vec{u}



"searchlight"

output: pulses w/ rep rate $1/\text{orbit period}$

Antennas - radiation from collections of charges

2 assumptions:

- 1) localized charge/current distribution (size d)
- 2) source oscillates w/ frequency ω ($\lambda = 2\pi c/\omega$)

Observation distance - r

Must make approximations to do calculations: order these dimensions

Scaling analysis:

$$\vec{E} = e \left[\frac{(\hat{R} - \beta)(1 - \beta^2)}{K^3 R^2} + \frac{\hat{R} \times ((\hat{R} - \beta) \times \vec{a})}{c^2 K^3 R} \right]$$

$$\vec{B} = \hat{R} \times \vec{E}, \quad \vec{S} = \frac{c}{4\pi} (\vec{E} \times \vec{H})$$

compare relative magnitude of terms.

K^3 common to all

single charge: oscillates over dist d at freq. $\nu =$

$$u \sim d \cdot \nu \quad \beta \sim d \nu / c \sim d / \lambda \quad a/c^2 \sim d / \lambda^2$$

since $\beta < 1$, single charge must have $d < \lambda$

example: $\nu \sim 1 \text{ MHz}$ (AM radio)

$$\rightarrow \lambda \sim 300 \text{ m}$$

$$d \sim 1 \text{ cm}$$

$$u \sim d \nu \sim 10^{-2} \text{ m} \cdot 10^6 \text{ s}^{-1} \sim 10^4 \text{ m/s}$$

$$\rightarrow \beta \sim d / \lambda \sim 3 \times 10^{-5}$$

$$\left. \begin{array}{l} 100 \text{ MHz FM} \\ 3 \text{ m} \end{array} \right\}$$

$$\left. \begin{array}{l} 10^6 \text{ m/s} \\ 3 \times 10^{-3} \end{array} \right\}$$

$$E \sim \frac{1}{R^2} + \frac{a}{c^2 R} \sim \frac{1}{R^2} + \frac{d}{\lambda^2 R}$$

$$S \sim \frac{1}{R^4} + \frac{2d}{\lambda^2 R^3} + \frac{d^2}{\lambda^4 R^2} \sim 1, \quad \frac{2dR}{\lambda^2}, \quad \frac{d^2 R^2}{\lambda^4}$$

Antennas - radiation from collections of charges.

consider charge distribution that is:

- localized
- oscillating with frequency ω

→ 3 length scales

size: d observation distance: r

wavelength of radiation: $\lambda = 2\pi c/\omega$

Ordering of these scales → method for solution.

Quasistatic

low frequency → long wavelength

e.g. $\nu = 1 \text{ MHz}$ (AM radio)

$$\lambda = 300 \text{ m}$$

observe w/in lab: $r \lesssim 2 \text{ m}$

$r \ll \lambda$

suppose source is $d \approx 1 \text{ cm}$

Can we ignore retardation effects? yes

$$d \ll \lambda$$

if charges are moving at speed d , $u \approx d \cdot \gamma \approx 10^{-2} \cdot 10^6$

$$u \ll c$$

$$= 10^4 \text{ m/s}$$

$$\beta \ll 1$$

$$a \approx d \cdot \gamma^2$$

$$K \equiv 1 - \vec{R} \cdot \vec{\beta} \approx 1$$

$$E \approx e \left[\frac{(\vec{R} - \vec{\beta})(1 - \beta^2)}{K^3 R^2} + \frac{\vec{R} \times ((\vec{R} - \vec{\beta}) \times \vec{a})}{c^2 K^3 R} \right]$$

$$\approx e \left[\frac{1}{R^2} + \frac{d \cdot \gamma^2}{c^2 R} \right] \approx e \left[\frac{1}{R^2} + \frac{d}{\lambda^2 R} \right]$$

$$\frac{u}{c} \approx \frac{d \cdot \gamma}{c} \approx \frac{d}{\lambda}$$

Now we are left with static fields (e.g. coulomb) that are slowly varying.

$$\rightarrow \Phi(\vec{r}, t) \approx \int \frac{\rho(\vec{r}', t)}{|\vec{r} - \vec{r}'|} d^3r'$$

if $r \gg d$ we can do a multipole expansion
e.g. calc. dipole, quadrupole moment.
fields $\sim 1/r^2$ or faster.

Radiation Fields:

if $r \gg \lambda$, static fields drop off quickly
 \rightarrow radiation

if source is small, do not ignore retardation
w/in source.

larger sources $d \sim \lambda$: superpositions