

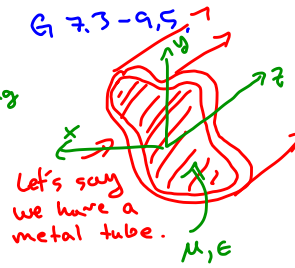
Readin Today: G 9.5

Next: Review G 7.3-9.5

Assume perfectly conducting walls,

→ Just inside the walls

$$\vec{E}_{||} = 0; \quad \vec{B}_{\perp} = 0$$



Maxwell's eqns inside waveguide.

$$\textcircled{1} \nabla \cdot \vec{E} = \rho \quad \textcircled{2} \nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$$

$$\textcircled{3} \nabla \cdot \vec{B} = 0 \quad \textcircled{4} \nabla \times \vec{B} = \mu \epsilon \frac{\partial \vec{E}}{\partial t}$$

If propagation is in z-direction then

$$\vec{E} = \vec{E}_0 e^{i(kz - \omega t)} \quad \vec{B} = \vec{B}_0 e^{i(kz - \omega t)}$$

\vec{E}_0, \vec{B}_0 depend only on x, y

$$\textcircled{3} \hat{z}: \frac{\partial E_z}{\partial x} - \frac{\partial E_x}{\partial y} = +i\omega B_z \quad \textcircled{4} \hat{z}: \frac{\partial B_y}{\partial x} - \frac{\partial B_x}{\partial y} = -\frac{i\omega}{v^2} E_z$$

$$\hat{y}: \frac{\partial E_z}{\partial y} - \frac{\partial E_y}{\partial z} = +i\omega B_x \quad \hat{x}: \frac{\partial B_z}{\partial y} - ik B_y = -\frac{i\omega}{v^2} E_x$$

$$\hat{y}: \frac{\partial E_x}{\partial z} - \frac{\partial E_z}{\partial x} = +i\omega B_y \quad \hat{y}: ik B_x - \frac{\partial B_z}{\partial x} = -\frac{i\omega}{v^2} E_y$$

$$ik E_x - \frac{\partial E_z}{\partial x} = \frac{i\omega}{ik} \left[\frac{\partial B_z}{\partial y} + \frac{i\omega}{v^2} E_x \right]$$

$$\Rightarrow -k^2 E_x - ik \frac{\partial E_z}{\partial x} = i\omega \frac{\partial B_z}{\partial y} - \frac{\omega^2}{v^2} E_x$$

$$E_x = \frac{i}{\omega^2/v^2 - k^2} \left(k \frac{\partial E_z}{\partial x} + \omega \frac{\partial B_z}{\partial y} \right)$$

Also you can get

$$E_y = \frac{i}{\omega^2/v^2 - k^2} \left(k \frac{\partial E_z}{\partial y} - \omega \frac{\partial B_z}{\partial x} \right)$$

$$B_x = \frac{i}{\omega^2/v^2 - k^2} \left(k \frac{\partial B_z}{\partial x} - \frac{\omega}{v^2} \frac{\partial E_z}{\partial y} \right)$$

$$B_y = \frac{i}{\omega^2/v^2 - k^2} \left(k \frac{\partial B_z}{\partial y} + \frac{\omega}{v^2} \frac{\partial E_z}{\partial x} \right)$$

$$E_x = \frac{i}{\omega^2/\mu^2 - k^2} \left(k \frac{\partial E_z}{\partial x} + \omega \frac{\partial B_z}{\partial y} \right)$$

Also you can get

$$E_y = \frac{i}{\omega^2/\mu^2 - k^2} \left(k \frac{\partial E_z}{\partial y} - \omega \frac{\partial B_z}{\partial x} \right)$$

$$B_x = \frac{i}{\omega^2/\mu^2 - k^2} \left(k \frac{\partial B_z}{\partial x} - \frac{\omega}{\nu^2} \frac{\partial E_z}{\partial y} \right)$$

$$B_y = \frac{i}{\omega^2/\mu^2 - k^2} \left(k \frac{\partial B_z}{\partial y} + \frac{\omega}{\nu^2} \frac{\partial E_z}{\partial x} \right)$$

Let's define $\vec{E}_T = E_x \hat{x} + E_y \hat{y} \Rightarrow \vec{E} = \vec{E}_T + E_z \hat{z}$
 $\vec{B}_T = B_x \hat{x} + B_y \hat{y} \Rightarrow \vec{B} = \vec{B}_T + B_z \hat{z}$

$$\vec{E}_T = \frac{i}{\omega^2/\mu^2 - k^2} \left[k \left[\frac{\partial E_z}{\partial x} \hat{x} + \frac{\partial E_z}{\partial y} \hat{y} \right] + \omega \left[\frac{\partial B_z}{\partial y} \hat{x} - \frac{\partial B_z}{\partial x} \hat{y} \right] \right]$$

Let's define an operator $\vec{\nabla}_T = \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial y} \hat{y}$
 $\vec{\nabla} = \vec{\nabla}_T + \frac{\partial}{\partial z} \hat{z}$

Show $\vec{\nabla}_T \cdot \vec{\nabla}_T f = \nabla_T^2 f$ $\{ \nabla_T^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \}$
 $\vec{\nabla}_T \cdot (\vec{\nabla}_T \times \vec{v}) = 0$

Rewriting \vec{E}_T and \vec{B}_T

$$\vec{E}_T = \frac{i}{\omega^2/\mu^2 - k^2} \left[k \vec{\nabla}_T E_z + \omega \vec{\nabla}_T \times \vec{B}_z \right]$$

$$\vec{B}_T = \frac{i}{\omega^2/\mu^2 - k^2} \left[k \vec{\nabla}_T B_z - \frac{\omega}{\nu^2} \vec{\nabla}_T \times \vec{E}_z \right]$$

Max (1): $\vec{\nabla} \cdot \vec{E} = 0$ $\vec{\nabla}_T \cdot \vec{E}_T + ik E_z = 0$

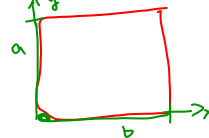
[Find a diff eq. for E_z using and

$$\vec{\nabla}_T \left[\frac{i}{\omega^2/\mu^2 - k^2} \left(k \vec{\nabla}_T E_z + \omega \vec{\nabla}_T \times \vec{B}_z \right) \right] + ik E_z = 0$$

$$\nabla_T^2 E_z + (\omega^2/\mu^2 - k^2) E_z = 0$$

$$\frac{\partial^2 E_z}{\partial x^2} + \frac{\partial^2 E_z}{\partial y^2} + (\omega^2/\mu^2 - k^2) E_z = 0$$

Same for B_z $\frac{\partial^2 B_z}{\partial x^2} + \frac{\partial^2 B_z}{\partial y^2} + (\omega^2/\mu^2 - k^2) B_z = 0$



Aside: TE (transverse electric)

$$\Rightarrow E_z = 0$$

TM (" magnetic)

$$\Rightarrow B_z = 0$$

TEM $\Rightarrow E_z = B_z = 0$.

The book treats the TE case.
 Lets start the TM case.

Separation of variables!

Assume $E_z = X(x)Y(y)$

$$\frac{\partial^2}{\partial x^2} X Y + \frac{\partial^2}{\partial y^2} X Y + (\frac{\omega^2}{v^2} - k^2) X Y = 0$$

$$Y \frac{d^2 X}{dx^2} + X \frac{d^2 Y}{dy^2} + (\frac{\omega^2}{v^2} - k^2) X Y = 0$$

$$\frac{1}{X} \frac{d^2 X}{dx^2} + \frac{1}{Y} \frac{d^2 Y}{dy^2} + (\frac{\omega^2}{v^2} - k^2) = 0$$

$$\Rightarrow \frac{1}{X} \frac{d^2 X}{dx^2} = -k_x^2 ; \frac{1}{Y} \frac{d^2 Y}{dy^2} = -k_y^2$$

$$X = A_1 \cos(k_x x) + A_2 \sin(k_x x) \quad Y = B_1 \cos(k_y y) + B_2 \sin(k_y y)$$

$$\vec{E}_{||} = 0 \Rightarrow E_z = 0 \text{ on the boundaries.}$$

$$\Rightarrow A_1 = B_1 = 0$$

$$X = A_2 \sin(k_x x) \quad Y = B_2 \sin(k_y y)$$

$$\frac{\omega^2}{v^2} - k^2 - k_x^2 - k_y^2 = 0$$

$$k_x^2 + k_y^2 + k^2 = \frac{\omega^2}{v^2}$$

$$k_x = \frac{n\pi}{b} \quad k_y = \frac{m\pi}{a} \quad \{n, m \text{ are integers}\}$$

Because there's an upper limit to $k_x^2 + k_y^2 < \frac{\omega^2}{v^2}$ then there is a finite # of propagating modes.

