

2 - 1 - 08

Note Title

1/31/2008

$$\Delta p \Delta x \geq \frac{\hbar}{2}$$

$\sigma_p \equiv \Delta x$ σ is more common in statistics.
Remember σ_p is the standard deviation

Lower Limit $\Delta x \Delta p = \hbar/2$

For QHO, total energy

$$E = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 x^2$$

must, therefore, be at least

$$\frac{(\Delta p)^2}{2m} + \frac{1}{2} m \omega^2 (\Delta x)^2$$

$\Delta p =$ momentum uncertainty

$\Delta x =$ position uncertainty

Using $\Delta p \Delta x = \hbar/2$ we have

$$E = \frac{1}{2m} \left(\frac{\hbar}{2\Delta x} \right)^2 + \frac{1}{2} m \omega^2 \Delta x^2$$

$$\frac{\hbar^2}{8m} (\Delta x)^{-2} + \frac{1}{2} m \omega^2 \Delta x^2$$

minimize this w.r.t. Δx

$$\frac{dE}{d(\Delta x)} \rightarrow 0$$

$$-\frac{\hbar}{4m} (\Delta x)^{-3} + m\omega^2 \Delta x = 0$$

$$\frac{\hbar}{4m} \frac{1}{(\Delta x)^3} = m\omega^2 \Delta x$$

$$\Delta x^4 = \frac{\hbar^2}{4m^2\omega^2}$$

$$\Delta x = \sqrt{\frac{\hbar}{2m\omega}}$$

insert this into

$$E = \frac{1}{2m} \left(\frac{\hbar}{2\Delta x} \right)^2 + \frac{1}{2} m\omega^2 \Delta x^2$$

to get

$$E_0 = \frac{\hbar^2}{8m} \frac{2m\omega}{\hbar} + \frac{1}{2} m\omega^2 \frac{\hbar}{2m\omega}$$

$$= \frac{\hbar\omega}{4} + \frac{\hbar\omega}{4} = \frac{\hbar\omega}{2}$$

This agrees with what we got before using a and a^\dagger .

The energy of a system governed by H.O. cannot be zero. So-called zero point vibration energy

E.g. specific heat of ideal gas (Einstein/Planck)

$$E = \frac{\hbar\omega}{e^{\hbar\omega/kT} - 1} + \frac{\hbar\omega}{2}$$

$$\text{as } T \rightarrow 0 \quad E \rightarrow \frac{\hbar\omega}{2}$$

E.g. keeps Helium 4 from solidifying as $T \rightarrow 0$.

E.g. Casimir Effect

$$\psi(x) = \sqrt{\frac{2}{a}} \sin\left(\frac{\pi x}{a}\right)$$

$$\langle p \rangle = \frac{2}{a} \int_0^a \sin\left(\frac{\pi x}{a}\right) (-i\hbar \frac{d}{dx}) \left(\sin\left(\frac{\pi x}{a}\right)\right) dx$$

$$= -i\hbar \frac{2}{a} \frac{2}{a} \int_0^a \sin\left(\frac{\pi x}{a}\right) \cos\left(\frac{\pi x}{a}\right) dx$$

$$\langle p^2 \rangle = + \hbar^2 \left(\frac{\pi}{a}\right)^2 \frac{2}{a} \int_0^a \sin^2\left(\frac{\pi x}{a}\right) dx$$

$$= 1$$

$$\sigma_p^2 = \langle p^2 \rangle - \langle p \rangle^2 = \hbar^2 \left(\frac{\pi}{a}\right)^2$$

$$\langle x \rangle = \frac{2}{a} \int_0^a x \sin^2\left(\frac{\pi x}{a}\right) dx$$

$$\frac{\pi x}{a} = y \quad dx = \frac{a}{\pi} dy$$

$$\frac{2}{a} \left(\frac{a}{\pi}\right)^2 \int_0^{\pi} y \sin^2(y) dy$$

$$\frac{2}{a} \left(\frac{a}{\pi}\right)^2 \left[\frac{y^2}{4} - \frac{1}{4} y \sin(2y) - \frac{1}{8} \cos(2y) \right]_0^{\pi}$$

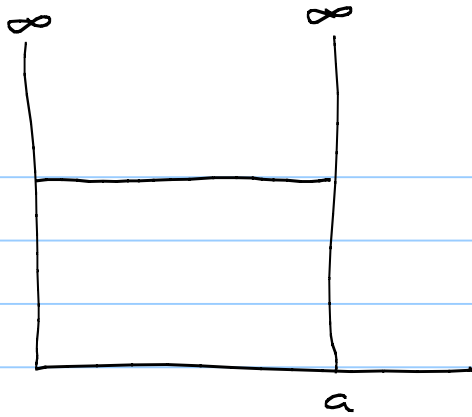
$$\frac{2}{a} \left(\frac{a}{\pi}\right)^2 \left[\left\{ \frac{\pi^2}{4} - \frac{1}{8} \right\} + \frac{1}{8} \right] = \frac{2}{a} \frac{a^2}{4\pi^2} = \frac{a}{2\pi^2}$$

$$\begin{aligned}\langle x^2 \rangle &= \frac{2}{a^2} \int_0^a x^2 \sin^2\left(\frac{\pi x}{a}\right) dx \\ &= \frac{1}{6} a^2 \left(2 - \frac{3}{\pi^2}\right)\end{aligned}$$

$$\begin{aligned}\sigma_x^2 &= \langle x^2 \rangle - \langle x \rangle^2 \\ &= \frac{1}{6} a^2 \left(2 - \frac{3}{\pi^2}\right) - \frac{a^2}{2^2} \\ &= a^2 \left[\frac{1}{3} - \frac{1}{4} - \frac{1}{2\pi^2} \right] \\ &= a^2 \left(\frac{1}{12} - \frac{1}{2\pi^2} \right)\end{aligned}$$

$$\begin{aligned}\sigma_x \sigma_p &= \hbar \left(\frac{\pi}{a}\right) a \sqrt{\left(\frac{1}{12} - \frac{1}{2\pi^2}\right)} \\ &= \hbar \sqrt{\frac{\pi^2 - 6}{12}} \approx \hbar 1.4\end{aligned}$$

So certainly $\sigma_x \sigma_p > \frac{\hbar}{2}$



$$f(x, 0) = \begin{cases} \frac{1}{\sqrt{a}} & 0 \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$

$$\frac{1}{\sqrt{a}} = \sum_{n=1}^{\infty} C_n \sqrt{\frac{2}{a}} \sin\left(\frac{2n\pi x}{a}\right)$$

$$C_n = \sqrt{\frac{1}{a}} \sqrt{\frac{2}{a}} \int_0^a \sin\left(\frac{2n\pi x}{a}\right) dx$$

$$= -\sqrt{\frac{1}{a}} \sqrt{\frac{2}{a}} \left[\frac{a}{2n\pi} \cos\left(\frac{2n\pi x}{a}\right) \right]_0^a$$

$$= -\sqrt{2} \frac{1}{2n\pi} \left[\cos(n\pi) - 1 \right]$$

$$C_n = \begin{cases} \sqrt{2} \frac{2}{n\pi} \frac{1}{2} & n \text{ odd} \\ 0 & n \text{ even} \end{cases}$$

(NB)

$$\sum_{n=1}^{\infty} |C_n|^2 = \frac{a}{\pi^2} \sum_{n=1,3,\dots}^{\infty} \frac{1}{n^2}$$

$$= \frac{a}{\pi^2} \left[1 + \frac{1}{9} + \frac{1}{25} + \dots \right] = 1$$

Proof $\sum_{n=1}^{\infty} \frac{1}{n^2} = \zeta(2) = \frac{\pi^2}{6}$

$$S \equiv 1 + \frac{1}{9} + \frac{1}{25} + \frac{1}{49} \dots$$

$$= \left(1 + \frac{1}{4} + \frac{1}{9} + \frac{1}{16} + \dots \right) - \left(\frac{1}{4} + \frac{1}{16} + \frac{1}{36} \dots \right)$$

$\underbrace{\hspace{15em}}_{\{2\}}$

$$= \sum_{n=1}^{\infty} \frac{1}{(2n)^2} = \frac{1}{4} \sum_{n=1}^{\infty} \frac{1}{n^2}$$

$$= \frac{1}{4} \zeta(2)$$

$$S = \zeta(2) - \frac{1}{4} \zeta(2) = \frac{3}{4} \zeta(2)$$

$$= \frac{3}{4} \frac{\pi^2}{6} = \frac{\pi^2}{8} \quad \text{🚩}$$

Back to the story

initial state uniformly distributed on $[0, a]$

$$\bar{\psi}(x, 0) = \sum_{n=1}^{\infty} c_n \sqrt{\frac{2}{a}} \sin\left(\frac{n\pi x}{a}\right)$$

$$c_n = \begin{cases} \frac{2}{\pi} \frac{1}{n} & \text{for } n \text{ odd} \\ 0 & \text{for } n \text{ even} \end{cases}$$