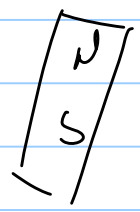


ϵ is small

$$\vec{B} = B_x(x, y, z) \hat{x} + B_y(x, y, z) \hat{y} + B_z(x, y, z) \hat{z}$$

4 $d\vec{r}$'s



$$d\vec{F} = I d\vec{\ell} \times \vec{B}$$

$$= I \left\{ \begin{array}{l} \text{lower} \\ dy \hat{y} \times \vec{B}(0, y, 0) + dz \hat{z} \times \vec{B}(0, \epsilon, z) \end{array} \right.$$

$$\left. \begin{array}{l} \text{upper} \\ - dy \hat{y} \times \vec{B}(0, y, \epsilon) - dz \hat{z} \times \vec{B}(0, 0, z) \end{array} \right\}$$

+ $\frac{1}{2}$ limits different

$$d\vec{F} = I \left\{ - dy \hat{y} \times \left[\vec{B}(0, y, \epsilon) - \vec{B}(0, y, 0) \right] + dz \hat{z} \times \left[\vec{B}(0, \epsilon, z) - \vec{B}(0, 0, z) \right] \right.$$

$$\left. \begin{array}{l} \epsilon \frac{\partial \vec{B}}{\partial z} \Big|_{0, y, 0} \qquad \qquad \qquad \epsilon \frac{\partial \vec{B}}{\partial y} \Big|_{0, 0, z} \end{array} \right.$$

Taylor series

$$\vec{B}(0, \epsilon, z) = \vec{B}(0, 0, z) + \epsilon \frac{\partial \vec{B}}{\partial y} \Big|_{0, 0, z} + \dots$$

$$\int_0^\epsilon dy \frac{\partial \vec{B}}{\partial z} \Big|_{0, y, 0} \approx \epsilon \frac{\partial \vec{B}}{\partial z} \Big|_{0, 0, 0}$$

constant $\frac{1}{2}$ same derivative as at $0, 0, 0$

$$\vec{F} = I \epsilon^2 \left[\hat{z} \times \frac{\partial \vec{B}}{\partial y} - \hat{y} \frac{\partial \vec{B}}{\partial z} \right]$$

↑
m magnetic dipole mom

$$\vec{F} = m \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 0 & 1 \\ \frac{\partial B_x}{\partial y} & \frac{\partial B_y}{\partial y} & \frac{\partial B_z}{\partial y} \end{vmatrix} - m \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ 0 & 1 & 0 \\ \frac{\partial B_x}{\partial z} & \frac{\partial B_y}{\partial z} & \frac{\partial B_z}{\partial z} \end{vmatrix}$$

$$= m \left\{ \hat{y} \frac{\partial B_x}{\partial y} - \hat{x} \frac{\partial B_y}{\partial y} - \hat{x} \frac{\partial B_z}{\partial z} + \hat{z} \frac{\partial B_x}{\partial z} \right\}$$

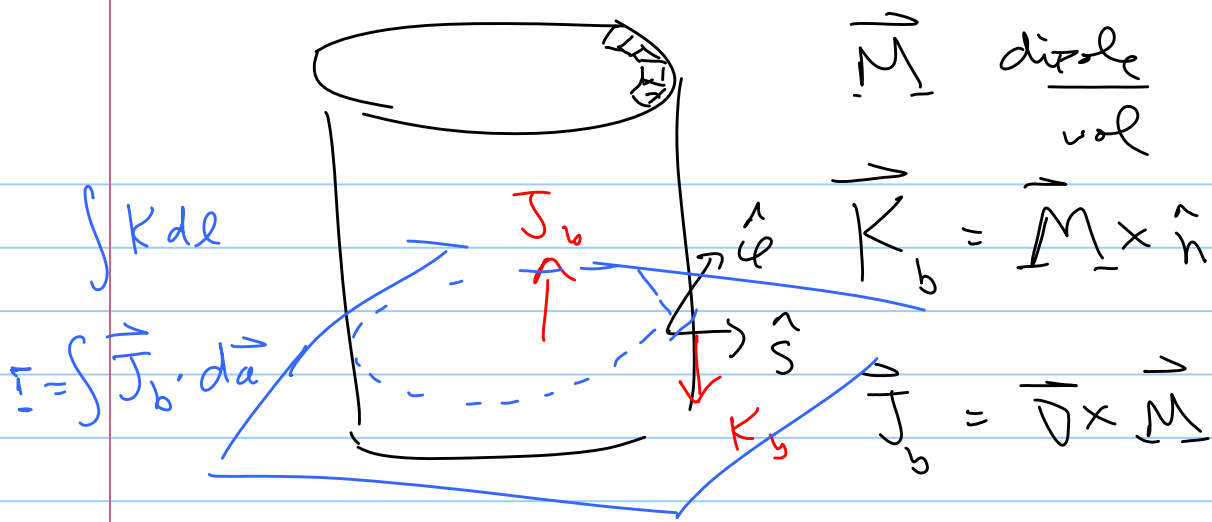
use $\vec{\nabla} \cdot \vec{B} = 0 = \frac{\partial B_x}{\partial x} + \frac{\partial B_y}{\partial y} + \frac{\partial B_z}{\partial z} \Rightarrow \frac{\partial B_x}{\partial x} = - \frac{\partial B_y}{\partial y} - \frac{\partial B_z}{\partial z}$

$$\vec{F} = m \left\{ \hat{x} \frac{\partial B_x}{\partial x} + \hat{y} \frac{\partial B_x}{\partial y} + \hat{z} \frac{\partial B_x}{\partial z} \right\}$$

But $\vec{m} \cdot \vec{B} = m B_x$ since $\vec{m} = m \hat{x}$

So $\vec{\nabla} (\vec{m} \cdot \vec{B}) = m \vec{\nabla} B_x = m \left(\frac{\partial B_x}{\partial x} \hat{x} + \frac{\partial B_x}{\partial y} \hat{y} + \frac{\partial B_x}{\partial z} \hat{z} \right)$
 $= \vec{F}$

$$\vec{F} = \vec{\nabla} (\vec{m} \cdot \vec{B})$$



$$\vec{M} = k s^2 \hat{\phi}$$

↑ radial direction

Prin: Find K_b & $J_b \Rightarrow$ Bist S or Amps

Method:

$$\vec{K}_b = \vec{M} \times \hat{n}; \hat{\phi} \times \hat{s}$$

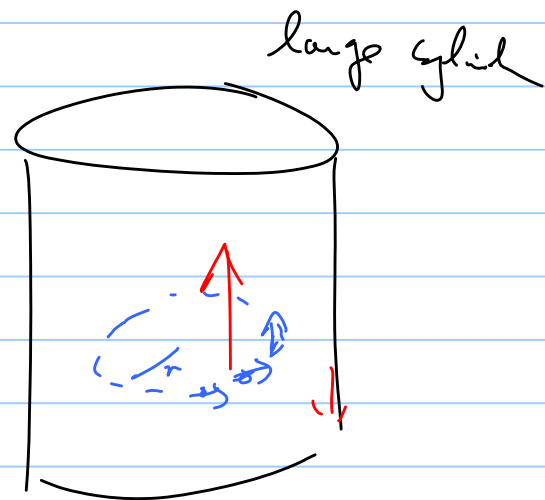
$$\vec{J}_b = \vec{\nabla} \times \vec{M}$$

B from Amp's Law

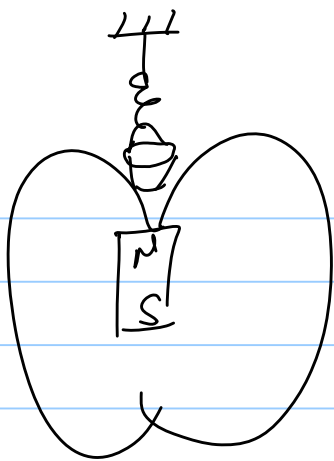
$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 I_{\text{enc}}$$

$$\Downarrow$$

$$B 2\pi r = \int \vec{J} \cdot d\vec{a}$$



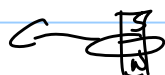
Real materials



- strongly attracted: Ferromag

- weak attraction: paramag

- repulsion: diamag
 Faradays law



mag statics

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J} = \mu_0 \vec{J}_{\text{free}} + \mu_0 \vec{J}_{\text{bound}}$$

$$\parallel \mu_0 \vec{\nabla} \times \vec{M}$$

$$\vec{\nabla} \times \left(\underbrace{\vec{B}}_{\mu_0} - \underbrace{\mu_0 \vec{M}}_{\mu_0} \right) = \vec{J}_f$$

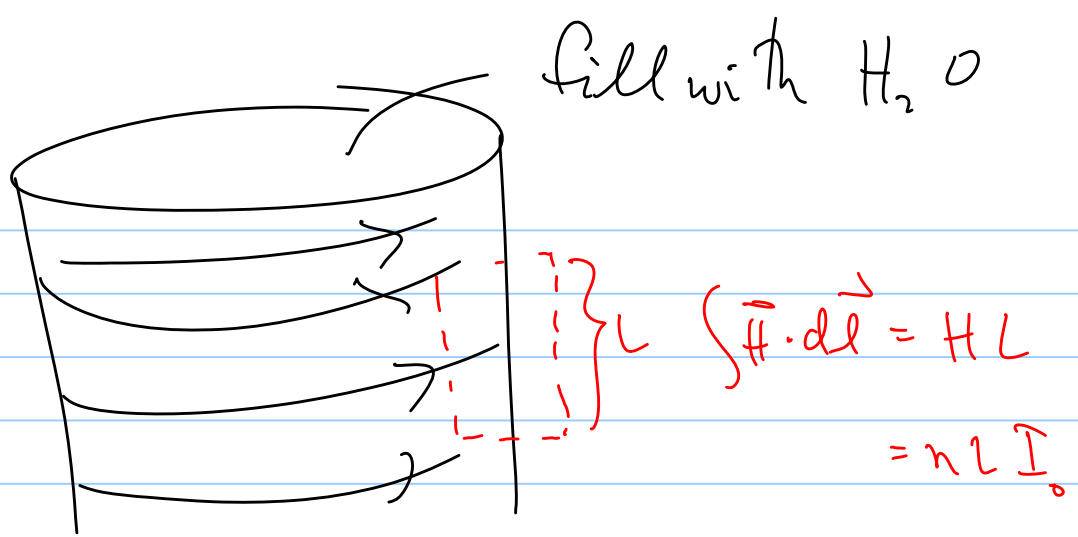
\vec{H}

$$\vec{\nabla} \times \vec{H} = \vec{J}_f \quad \oint \vec{H} \cdot d\vec{e} = I_f$$

\uparrow
 $\cdot da$

\uparrow
 $\cdot da$

Stokes Th



find $H = nI$

Assume linear material

$\vec{M} \propto \vec{H}$

paramag pos
 diam neg

$\vec{M} = \chi_m \vec{H}$

$$\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} = \frac{\vec{B}}{\mu_0} - \chi_m \vec{H}$$

$$\mu_0(1 + \chi_m) \vec{H} = \vec{B}$$

$$\mu \vec{H} = \vec{B}$$