

Schrödinger eqn calc of  $\chi_{-}$

$\hat{H}_0$  is hamiltonian for atom  
 $\hat{V}(t) = -\hat{\vec{\mu}} \cdot \vec{E}(t)$  is perturbing interaction  
 $\hat{\vec{\mu}} = -e\vec{r}$  dipole operator  
 $\vec{E}(t) = \sum_p \vec{E}(w_p) e^{-i w_p t}$   $p$  is  $\pm 1$

unperturbed eigenstates of atom:

$\Psi_n(\vec{r}, t) = u_n(\vec{r}) e^{-i w_n t}$   
 so that  $\hat{H}_0 u_n = E_n u_n$   
 w/ orthonormality:  
 $\int u_m^* u_n d^3r = \delta_{mn}$

Time-dependent perturbation theory.

Full Hamiltonian is

$\hat{H} = \hat{H}_0 + \lambda \hat{V}(t)$   $0 < \lambda < 1$

if  $\Psi(\vec{r}, t)$  is a solution to the full  $\hat{H}$ :  $\hat{H}\Psi = i\hbar \frac{\partial \Psi}{\partial t}$

We can express this solution as an expansion:

$\Psi(\vec{r}, t) = \Psi^{(0)}(\vec{r}, t) + \lambda \Psi^{(1)}(\vec{r}, t) + \lambda^2 \Psi^{(2)}(\vec{r}, t) + \dots$

collect terms of equal order in  $\lambda$

0<sup>th</sup> order:  $i\hbar \frac{\partial \Psi^{(0)}}{\partial t} = \hat{H}_0 \Psi^{(0)}$  unperturbed case

1<sup>st</sup> order:  $i\hbar \frac{\partial \Psi^{(1)}}{\partial t} = \hat{H}_0 \Psi^{(1)} + \hat{V} \Psi^{(0)}$

N<sup>th</sup> order:  $i\hbar \frac{\partial \Psi^{(N)}}{\partial t} = \hat{H}_0 \Psi^{(N)} + \hat{V} \Psi^{(N-1)}$

we'll assume the electron is initially in the ground state.

$\Psi(\vec{r}, t=0) = u_g(\vec{r})$   $\Psi^{(0)}(\vec{r}, t) = u_g(\vec{r}) e^{-i E_g t / \hbar}$

Now we assume that unpert. solutions forms a complete set.

we can represent any order  $N$  in that basis

$$\Psi^{(N)}(\vec{r}, t) = \sum_l a_l^{(N)}(t) U_l(\vec{r}) e^{-i\omega_l t} \quad \omega_l = E_l/\hbar$$

in first-order PT we're looking for  $\Psi^{(1)}$

general case

$$\rightarrow i\hbar \sum_l \dot{a}_l^{(1)} U_l(\vec{r}) e^{-i\omega_l t} = \sum_l a_l^{(0)} V U_l(\vec{r}) e^{-i\omega_l t}$$

1<sup>st</sup> order: if there is only the ground state at  $t=0$ ,

$$i\hbar \sum_l \dot{a}_l^{(1)} U_l(\vec{r}) e^{-i\omega_l t} = V U_0(\vec{r}) e^{-i\omega_0 t}$$

We want to extract an eqn. for the coeffs:

mult by  $U_m^*(\vec{r})$  and integrate. use orthogonality

$$\begin{aligned} \rightarrow i\hbar \sum_l \dot{a}_l^{(1)} \int U_m^* U_l e^{-i\omega_l t} d^3r &= i\hbar \dot{a}_m^{(1)} e^{-i\omega_m t} \\ &= \sum_l a_l^{(0)} \int U_m^* V U_l e^{-i\omega_l t} d^3r \equiv \sum_l a_l^{(0)} V_{ml} e^{-i\omega_l t} \end{aligned}$$

Finally, we have set of eqns:

$$\dot{a}_m^{(1)} = \frac{1}{i\hbar} \sum_l a_l^{(0)} V_{ml} e^{i\omega_{ml} t} \quad \omega_{ml} = \omega_m - \omega_l$$

This is actually a matrix equation. Suppose  $n$  levels

$$\begin{pmatrix} \dot{a}_1^{(1)} \\ \dot{a}_2^{(1)} \\ \dot{a}_n^{(1)} \end{pmatrix} = \frac{1}{i\hbar} \begin{pmatrix} V_{11} & V_{12} & V_{1n} \\ V_{21} & V_{22} & V_{2n} \\ V_{n1} & V_{n2} & V_{nn} \end{pmatrix} \begin{pmatrix} a_1^{(0)} \\ a_2^{(0)} \\ a_n^{(0)} \end{pmatrix}$$

insert  $e^{i\omega_{ml} t}$  terms on off diagonals.

we can evaluate  $a_m^{(1)}(t)$  by integrating:

$$a_m^{(1)}(t) = \frac{1}{i\hbar} \sum_l \int_{-\infty}^t dt' V_{ml}(t') a_l^{(0)}(t') e^{i\omega_{ml} t'}$$

Solutions:

1) system initially in ground state

$$2) V_{int}(t') = - \sum_p \vec{\mu}_{int} \cdot \vec{E}(\omega_p) e^{-i\omega_p t'}$$

$$\vec{\mu}_{int} = \int U_m^* \vec{\mu} U_g d^3r = \text{electric dipole transition moment.}$$

3) calc. 1<sup>st</sup> order term  $\rightarrow$  linear response,  $\chi^{(1)}$

$$a_m^{(1)}(t) = \frac{1}{\hbar} \sum_p \frac{\vec{\mu}_{mg} \cdot \vec{E}(\omega_p)}{\omega_{mg} - \omega_p} e^{i(\omega_{mg} - \omega_p)t}$$

4) calc. 2<sup>nd</sup> order term: involves a second sum over input fields

$\rightarrow \chi^{(2)}$

$$a_n^{(2)}(t) = \frac{1}{\hbar^2} \sum_{pq} \sum_m \frac{(\vec{\mu}_{nm} \cdot \vec{E}(\omega_q)) (\vec{\mu}_{mg} \cdot \vec{E}(\omega_p))}{(\omega_{ng} - \omega_p - \omega_q) (\omega_{ng} - \omega_p)} e^{i(\omega_{ng} - \omega_p - \omega_q)t}$$

add up  $\omega$ 's

$$V_{nm} a_m e^{i\omega_{nm}t} \\ \downarrow \quad \downarrow \quad \downarrow \\ -\omega_q + \omega_{mg} - \omega_p + \omega_{nm} = \omega_{ng} - \omega_p - \omega_q.$$

5) 3<sup>rd</sup> order ...

Note direct connection btw order of TDPT and  $\chi$  order, # photons

Calculation of suscept.

Linear

$$\text{recall classical calc: } \vec{P} = \chi^{(1)} \vec{E} = N_a \vec{P}$$

atomic density  $\cdot$  dipole moment.

$$QM: \langle \vec{P} \rangle = \langle \Psi | \vec{\mu} | \Psi \rangle \quad \text{expectation value.}$$

remember we are expanding  $\Psi$  in pert. series.

$\rightarrow$  expansion of  $\vec{P}$

$$\langle P^{(1)} \rangle = \langle \Psi^{(0)} | \vec{\mu} | \Psi^{(1)} \rangle + \langle \Psi^{(1)} | \vec{\mu} | \Psi^{(0)} \rangle$$

$$|\Psi^{(1)}\rangle = \sum_m a_m^{(1)}(t) |u_m\rangle e^{-i\omega_m t}$$