

Maxwell's eqns.

Note Title

6/19/2006

① $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon_0}$ ② $\nabla \cdot \vec{B} = 0$

③ $\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$ ④ $\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$

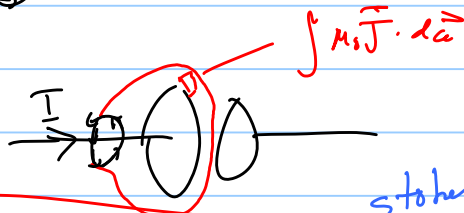
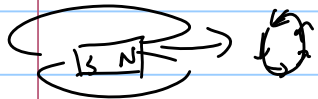


$\int \nabla \times \vec{B} \cdot d\vec{a} = \oint \vec{B} \cdot d\vec{l}$

$\int \nabla \times \vec{E} \cdot d\vec{a} = \oint \vec{E} \cdot d\vec{l} = -\frac{\partial}{\partial t} \int \vec{B} \cdot d\vec{a}$

$\text{EMF} \qquad \qquad \qquad \Phi_m$

$\int \mu_0 \vec{J} \cdot d\vec{a} = \mu_0 I$



stokes theorem

not s.s.

$\int \vec{B} \cdot d\vec{l} = \mu_0 I + \int \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t} \cdot d\vec{a}$

\uparrow
steady state

$\nabla \times \vec{B} = \mu_0 \vec{J} + \mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t}$

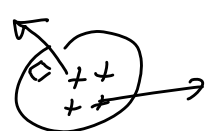
$\int \vec{B} \cdot d\vec{l}$

$\nabla \cdot (\nabla \times \vec{B}) = 0 = \mu_0 \nabla \cdot \vec{J} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \nabla \cdot \vec{E} \equiv 0$

Cartesian $\nabla \cdot \vec{C} = \frac{\partial}{\partial x} C_x + \frac{\partial}{\partial y} C_y + \frac{\partial}{\partial z} C_z = 0$

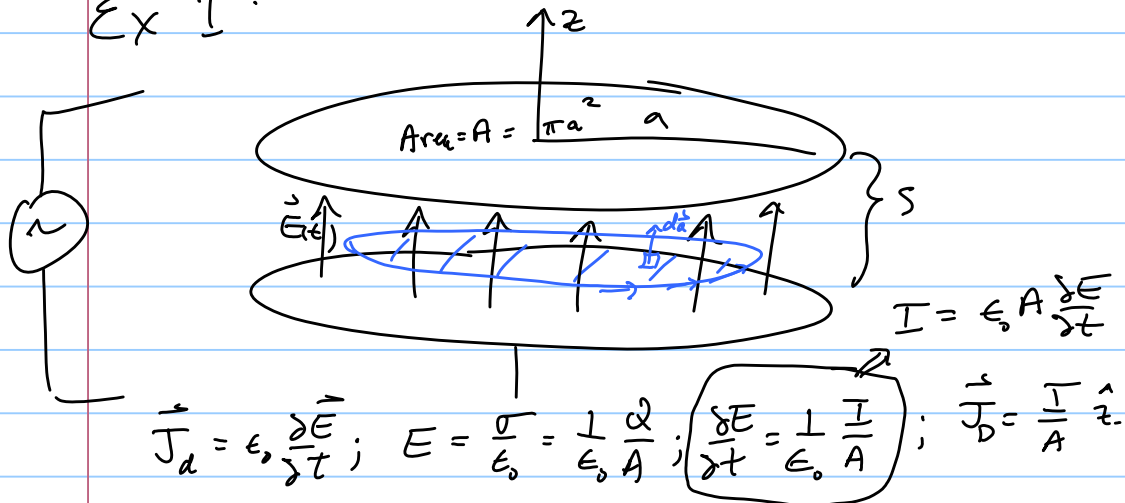
$$\boxed{\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}} \quad \begin{array}{l} \text{conservation of charge} \\ \frac{Q}{s} \text{ flowing out of volume} \end{array}$$

$$\int \nabla \cdot \vec{J} d\tau = \oint \vec{J} \cdot d\vec{a} = -\frac{dQ_{\text{enc}}}{dt} = -\frac{d}{dt} \int \rho d\tau = \int -\frac{\partial \rho}{\partial t} d\tau$$

$$\vec{J} = \rho \vec{v} \quad \frac{C}{m^3 s} \quad \frac{C}{s m^2}$$


- Maxwell's eqns which give $\vec{E} \neq \vec{B}$
- conservation of charge
- $\vec{F} = q\vec{E} + q\vec{v} \times \vec{B} = m\vec{a}$

Ex 1.

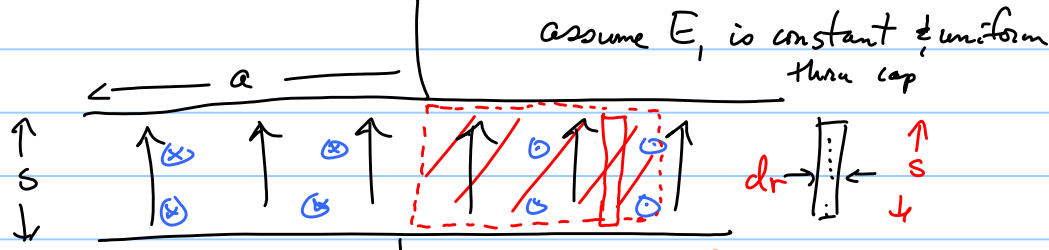


$$\int \vec{B} \cdot d\vec{l} = \mu_0 \int \vec{J}_d \cdot d\vec{a} = \mu_0 \frac{I}{\pi a^2} \pi r^2 \quad \boxed{\vec{B} = \frac{\mu_0 I}{2\pi} \frac{r}{a^2} \hat{\phi}}$$

Changing \vec{E} creates \vec{B} (like Faraday's law)

We have $I(t)$ so $B(t) \neq$ changing B must generate an E.

Approach is iterative soln.



$$\oint \vec{E} \cdot d\vec{l} = \oint (\vec{E}_1 + \vec{E}_2 + \dots) \cdot d\vec{l} = \oint \vec{E}_1 \cdot d\vec{l} + \oint \vec{E}_2 \cdot d\vec{l}$$

define $\boxed{\vec{E}(r=0) = \vec{E}_0 e^{i\omega t}} = \vec{E}_1 + \vec{E}_2$

$$E_2(r=0) = 0$$

$$\oint \vec{E} \cdot d\vec{l} = \oint \vec{E}_2 \cdot d\vec{l} = -E_2 s = -\frac{\partial \Phi_m}{\partial t} = -\frac{\partial}{\partial t} \int B_1(r) da$$

$$= -\frac{\partial}{\partial t} \int B_1(r) s dr = -\frac{\partial}{\partial t} \int_0^a \left(\frac{\mu_0 I}{2\pi} \frac{r}{a^2} \right) s dr$$

$$I = \epsilon_0 A \frac{\partial E}{\partial t} = \epsilon_0 A i \omega E_0 e^{i\omega t}$$

B has a phase shift

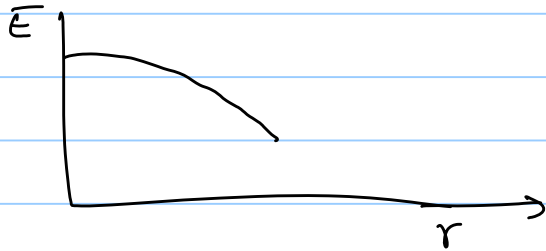
$$B_1 = \frac{\mu_0 r}{2\pi a^2} \epsilon_0 A i \omega E_0 e^{i\omega t} = \frac{i\omega}{2} \mu_0 \epsilon_0 r E_0 e^{i\omega t}$$

$$-E_2 = -\frac{d}{dt} \int \frac{i\omega}{2} \mu_0 \epsilon_0 r E_0 e^{i\omega t} dr$$

$$E_2 = -\frac{\omega^2 r^2}{4} \mu_0 \epsilon_0 E_0 e^{i\omega t}$$

$$c^2 = \frac{1}{\mu_0 \epsilon_0}$$

$$E = E_1 + E_2 = E_0 e^{i\omega t} \left(1 - \frac{1}{4} \frac{\omega^2 r^2}{c^2} \right)$$



Now the mag. field we cal. isn't right because the changing E used to get B was wrong. $B = B_1 + B_2$

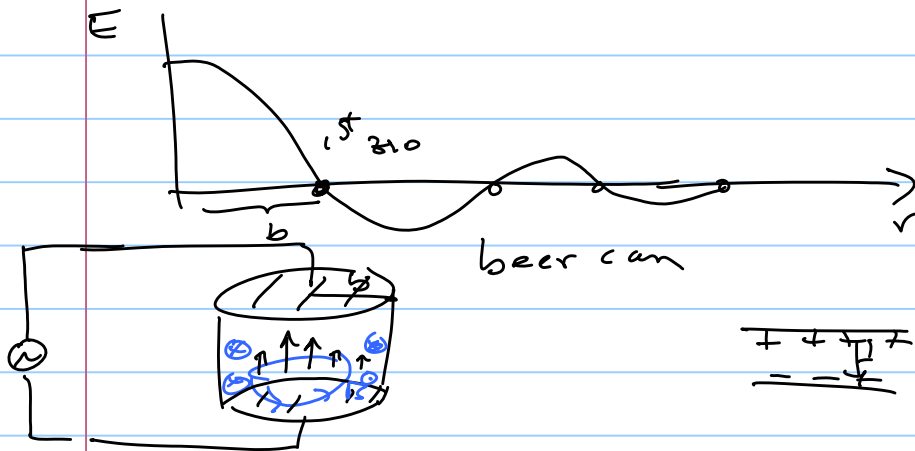
Where $B_1 = \frac{i\omega}{2} \frac{r}{c^2} E_0 e^{i\omega t}$ due to variations in E_1 ;

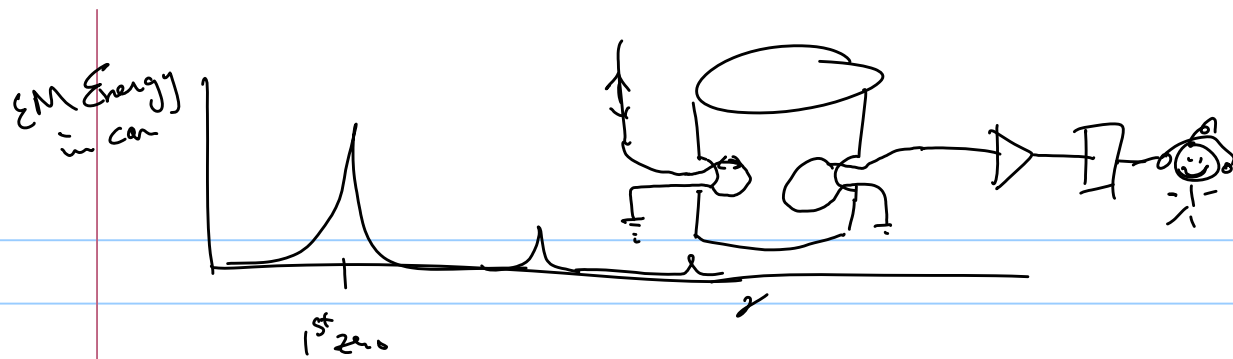
B_2 is due to variations in E_2 . Go back to 1st figure and redo calculation with just E_2 to get B_2

$$E = E_0 e^{i\omega t} \left[1 - \frac{1}{2^2} \left(\frac{\omega r}{c} \right)^2 + \frac{1}{2^2 4^2} \left(\frac{\omega r}{c} \right)^4 + \dots \right]$$

↑
dimensionless

$$J_0(x) = 1 - \frac{1}{(1!)^2} \left(\frac{x}{2} \right)^2 + \frac{1}{(2!)^2} \left(\frac{x}{2} \right)^4 - \frac{1}{(3!)^2} \left(\frac{x}{2} \right)^6 + \dots$$





Klystron

