

Closed box: blackbody cavity

3-D solution to wave eqn can be written as

$$E \sim A e^{ik_x x} e^{ik_y y} e^{ik_z z} e^{-i\omega t}$$

or

$$A e^{i(\vec{k} \cdot \vec{r} - \omega t)}$$

\vec{k} is wave vector

boundary condition for wave along z

$E_x = 0$ at z walls \rightarrow use \sin ,

$E_x = 0$ along x walls " "

$\frac{\partial E_x}{\partial y} = 0$ along y walls " \cos ;

$\rightarrow E_x(x, y, z) = A_y \sin k_x x \cos k_y y \sin k_z z$
other permutations work too:

$$A_x \cos k_x x \sin k_y y \sin k_z z$$

$$A_z \sin k_x x \sin k_y y \cos k_z z$$

into wave eqn,

$$\rightarrow k_x^2 + k_y^2 + k_z^2 = \frac{\omega^2}{c^2} = k^2$$

$$\text{with } k_x = \frac{l\pi}{L_x} \quad k_y = \frac{m\pi}{L_y} \quad k_z = \frac{n\pi}{L_z}$$

\rightarrow allowed frequencies

$$\omega_{lmn} = c \sqrt{\left(\frac{l\pi}{L_x}\right)^2 + \left(\frac{m\pi}{L_y}\right)^2 + \left(\frac{n\pi}{L_z}\right)^2}$$

$$\text{or } \nu_{lmn} = c \sqrt{\left(\frac{l}{2L_x}\right)^2 + \left(\frac{m}{2L_y}\right)^2 + \left(\frac{n}{2L_z}\right)^2}$$

At all permutations, there are 6 modes for each l, m, n

Now we assume that the cavity walls are held at a temperature T .

If the EM mode energy is in thermal equilibrium with the walls there is a probability that that mode will be excited: $P(E)$

Original assumption: $P(E) \propto e^{-E/kT}$ (Maxwell-Boltzmann)

→ avg. energy of each mode is

$$\langle E \rangle = \frac{\int_0^\infty E e^{-E/kT} dE}{\int_0^\infty e^{-E/kT} dE} = kT$$

this assumes a continuous range of allowed energies.
(will be relevant to the next part)

$\langle E \rangle = kT$ gives energy in each mode.

need # modes/volume of box/unit $d\omega = \rho_\omega$

ω = frequency of EM wave (quantized b/c of boundary conditions)

$\rho_\omega \langle E \rangle$ = spectral energy density = ρ_ω

so that total energy density is $\int_0^\infty \rho_\omega d\omega$

Counting # modes:



In k -space, modes are equally spaced along each axis

each combination of $k_x, k_y, k_z \rightarrow \vec{k}$

many \vec{k} 's that give same $|\vec{k}|$ or ω

i.e. count # modes in a volume of k -space.

$$\begin{aligned}
 N(\nu) &= \# \text{ modes w/ freq } \nu \text{ or less, } |\vec{k}| \leq \frac{2\pi\nu}{c} \\
 &= 2 \cdot \frac{1}{8} \frac{\frac{4}{3}\pi \left(\frac{2\pi\nu}{c}\right)^3}{\frac{\pi}{2L_x} \frac{\pi}{2L_y} \frac{\pi}{2L_z}} \rightarrow \text{vol. of sphere in } k\text{-space} \\
 &\quad \uparrow \quad \uparrow \quad \quad \quad \rightarrow \text{vol. of } k\text{-space occ. by} \\
 &\quad \text{polars} \quad \quad \quad \text{1 mode} \\
 &\quad \quad \quad \text{only 1st octant segment.}
 \end{aligned}$$

$$= \frac{8\pi\nu^3}{3c^3} V \quad V = L_x L_y L_z = \text{vol. of box}$$

$$\text{now } p_\nu = \frac{1}{V} \frac{dN}{d\nu} = \frac{8\pi\nu^2}{c^3}$$

As $\nu \uparrow$ many more possible modes.

\therefore energy density is

$$p_\nu = p_\nu \langle E \rangle = \frac{8\pi\nu^2}{c^3} kT \quad \text{Rayleigh - Jeans}$$

$$\text{but total energy density is } \int_0^\infty \frac{8\pi\nu^2}{c^3} kT d\nu = \infty$$

Planck: energy in each mode is quantized

$$E = n h\nu \quad \text{not continuous.}$$

$\therefore \langle E \rangle$ is a sum not integral

$$= \frac{\sum_n n h\nu e^{-n h\nu/kT}}{\sum_n e^{-n h\nu/kT}} = \frac{h\nu}{e^{h\nu/kT} - 1}$$

for $h\nu \rightarrow 0$ same as R-L-B

but as $h\nu \sim kT$ denominator in $\langle E \rangle$ gets large

$$\therefore \text{black body curve } p_\nu = \frac{8\pi\nu^2}{c^3} \frac{h\nu}{e^{h\nu/kT} - 1}$$

Comparing w/ expt $\rightarrow h = 6.62 \times 10^{-34} \text{ Js}$ Planck's constant



- total amt of radiation \uparrow with T (Stefan-Boltzmann)
 $E_{\text{tot}} \propto T^4$
- note that curve was diff't peak when measured in λ space.
- BB expr gives photon occupation / mode present in vacuum.

$$\langle \phi \rangle \equiv \frac{\langle E \rangle}{h\nu} = \frac{1}{e^{h\nu/kT} - 1}$$

very small at room temp $kT \sim 1/40 \text{ eV}$ for visible light $h\nu \sim 2 \text{ eV}$

Energy density of EM waves:

$$\rho = \left\langle \frac{1}{2} \epsilon E^2 \right\rangle + \left\langle \frac{1}{2} \mu H^2 \right\rangle$$

↪ time average

return to 1-D resonator:



- allowed λ or ν is quantized by walls (1st quantization)
but classically field strength is continuous.

- QED: Hamiltonian for field (= energy)
 $H_f = \frac{1}{2} \epsilon E_f^2 + \frac{1}{2} \mu H_f^2$ for a given cavity mode.

Solve by analogy w/ SHD:

$$H = \frac{1}{2} \frac{\hat{p}_x^2}{m} + \frac{1}{2} m \omega^2 \hat{x}^2$$

$$\begin{aligned} \hat{p}_x &\rightarrow E_\nu \\ \hat{x} &\rightarrow H_\nu \end{aligned}$$

→ discrete allowed energies, separated by $h\nu$

(2nd quantization) $\rho_\nu = (n + \frac{1}{2}) h\nu / V$

min energy for $n_\nu = 0 = \frac{1}{2} \frac{h\nu}{V} = \text{vac fluctuations.}$

Laser cavity can put many photons in same mode.

- photons are bosons, follow Bose-Einstein stats