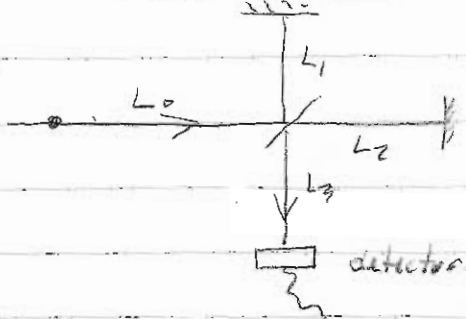


Michelson interferometer



Assume 50-50 beamsplit

$$\text{input } I_0 = \frac{c}{8\pi} E_0^2$$

intensity is split 50%

$$\rightarrow E_1 = \frac{1}{\sqrt{2}} E_0 = E_2$$

2 methods

1) optical path difference -

$$E_1 = E_0 e^{ik_0 L_0} \cdot \frac{1}{\sqrt{2}} e^{2ik_0 L_1} \cdot \frac{1}{\sqrt{2}} e^{ik_0 L_3} e^{-i\omega t}$$

$$= \frac{1}{2} E_0 \exp(ik_0 (L_0 + 2L_1 + L_3)) e^{-i\omega t}$$

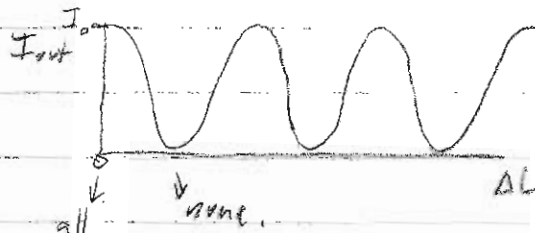
$$E_2 = \frac{1}{2} E_0 \exp(ik_0 (L_0 + 2L_2 + L_3)) e^{-i\omega t}$$

$$E_{\text{out}} = E_1 + E_2 = \frac{1}{2} E_0 e^{ik_0 (L_0 + L_3)} \left(e^{2ik_0 L_1} + e^{2ik_0 L_2} \right)$$

Note similarity to method when working with polarisation.

$$\rightarrow \Delta\phi = \delta = 2k_0 (L_1 - L_2)$$

$$I_{\text{out}} = \frac{1}{4} \frac{c}{8\pi} E_0^2 |1 + e^{i\delta}|^2 = \frac{1}{2} I_0 (1 + \cos 2k_0 (L_1 - L_2))$$



I_0 is indep of common phase

2) time delay

$$\text{path 1 } \tau_1 = (L_0 + 2L_1 + L_3)/c$$

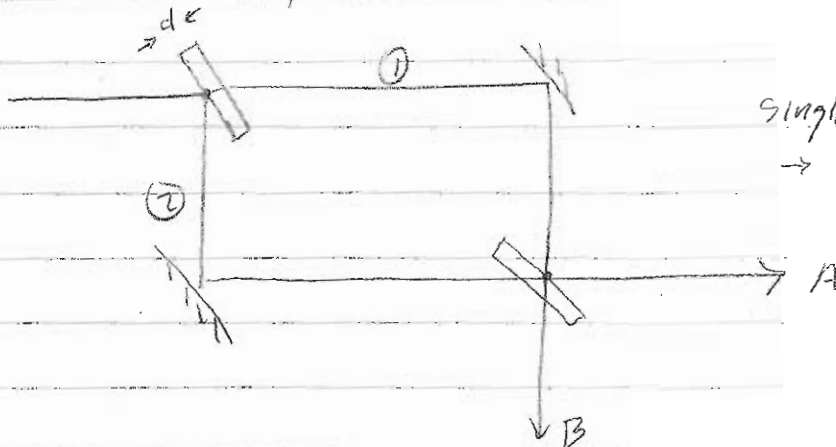
$$\text{path 2 } \tau_2 = (L_0 + 2L_2 + L_3)/c$$

$$\text{then } \Delta\phi = \frac{\omega}{c} (2L_1 - 2L_2) = k_0 (2L_1 - 2L_2) \text{ same}$$

Where does energy go? Out other port.

Mach-Zehnder interferometer: can measure both ports.

windows as beam splitters $|r|^2 \approx 4\%$



single pass thru

$$\rightarrow \phi_w = k_0 n d \cos 45^\circ$$

$$E_{1A} = E_0 \sqrt{1-R} (+i) e^{i\phi_w} e^{ik_0 L_1} \sqrt{R'} (L-1)$$

$$E_{2A} = E_0 \sqrt{R'} (L-1) e^{ik_0 L_2} e^{i\phi_w} \sqrt{1-R} (+i)$$

$$\rightarrow \delta = k_0 (L_1 - L_2) \text{ same amplitude.}$$

$$E_{1B} = E_0 \sqrt{1-R} (+i) e^{i\phi_w} e^{ik_0 L_1} \sqrt{1-R'} (L+1) e^{i\phi_w}$$

$$= E_0 (1-R) e^{i2\phi_w} e^{ik_0 L_1}$$

$$E_{2B} = E_0 \sqrt{R'} (L-1) e^{ik_0 L_2} e^{i2\phi_w} (+i) \sqrt{R'}$$

$$= E_0 R e^{i2\phi_w} (-e^{ik_0 L_2})$$

$$I_A = I_0 \cdot R(1-R) (1 + \cos k_0 (L_1 - L_2)) \quad \delta = \pi$$

$$I_A = 0$$

$$I_B = \frac{c}{8\pi} E_0^2 \left| (1-R) - R e^{i\pi(L_2-L_1)} \right|^2$$

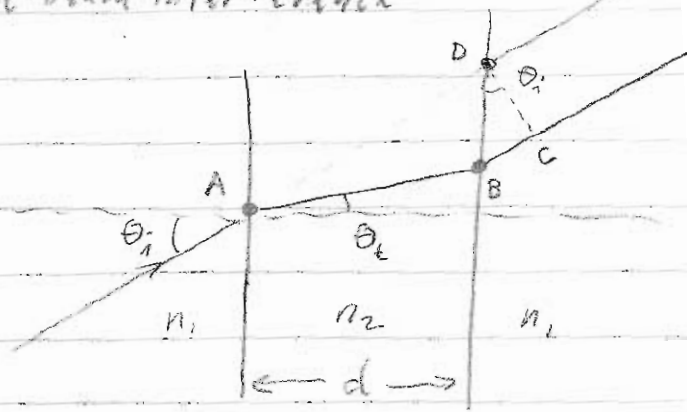
$$= I_0 \left((1-R)^2 + R^2 - (1-R)R (e^{i\pi(L_2-L_1)} + e^{-i\pi(L_2-L_1)}) \right)$$

$$= I_0 (1 - 2R + 2R^2 - 2(1-R)R \cos \delta) \quad \delta = \pi$$

$$I_B = I_0 (1 - 2(1-R)R (1 + \cos \delta))$$

$$I_{T2} = I_0$$

Double-beam interference



Suppose we introduce a tilted window into one arm of an interferometer. What is the phase shift?

simple case:



$$S_1 = d$$

$$S_2 = nd$$

$$\Delta S = (n-1)d$$

$$\Delta\phi = k_0(n-1)d$$

now tilt:

from diagram above:

try this: path in glass is $d/\cos\theta_t = \overline{AB}$

$$\rightarrow AS = \frac{n_2 d}{\cos\theta_t} = n_2 d \quad \text{X!}$$

this is wrong! We must be more careful.

geometric method: measure all the way out to line \perp to external ray

$$S_1 = n_1 \overline{AD} = n_1 d / \cos\theta_i$$

$$S_2 = n_2 \overline{AB} + \overline{BC} = n_2 d / \cos\theta_t + n_1 \overline{BD} \sin\theta_i$$

$$\overline{BD} = d \tan\theta_i - d \tan\theta_t$$

$$S_2 = \frac{n_2 d}{\cos\theta_t} + n_1 d (\tan\theta_i - \tan\theta_t) \sin\theta_i$$

$$n_1 \sin\theta_i = n_2 \sin\theta_t$$

$$S_2 = \frac{n_2 d}{\cos \theta_t} - \frac{n_2 d \sin^2 \theta_t}{\cos \theta_t} + n_1 d \sin \theta_i \tan \theta_i$$

$$= \frac{n_2 d \cos \theta_t + n_1 d \sin^2 \theta_i}{\cos \theta_i}$$

$$S_2 - S_1 = \frac{n_2 d \cos \theta_t + n_1 d \sin^2 \theta_i}{\cos \theta_i} - n_1 d \frac{1}{\cos \theta_i}$$

$$= \frac{n_2 d \cos \theta_t - n_1 d \cos \theta_i}{\cos \theta_i}$$

wave method:

$$E_1(y, z) = E_0 e^{i(k_{1y}y + k_{1z}z)}$$

$$\text{where } k_{1y} = k_0 n_1 \sin \theta_i$$

$$k_{1z} = k_0 n_1 \cos \theta_i$$

$$E_2(y, z) = E_0 e^{i(k_{2y}y + k_{2z}z)}$$

$$k_{2y} = k_0 n_2 \sin \theta_t$$

$$k_{2z} = k_0 n_2 \cos \theta_t$$

$$\text{but } k_{1y} = k_{2y}$$

evaluate waves at $z=d$

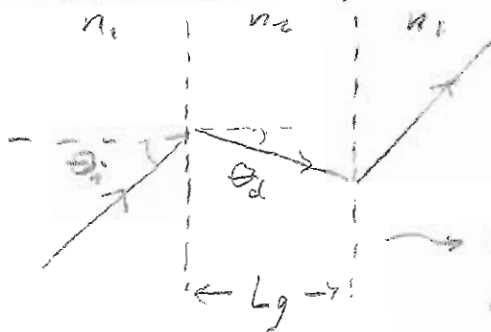
$$E_1 + E_2 = E_0 e^{-ik_{1y}y} \left[e^{ik_0 n_2 d \cos \theta_t} + e^{ik_0 n_1 d \cos \theta_i} \right]$$

$$\text{can see } \Delta \phi = k_0 d (n_2 \cos \theta_t - n_1 \cos \theta_i)$$

Note that this wave method can be extended to calculate the spectral phase of grating pairs and prism pairs:

Grating pairs

consider a pair of parallel transmission gratings:



$$n_2 \sin \theta_d = n_1 \sin \theta_i + \frac{m \lambda_0}{d}$$

grating spacing d
diffracted order m

$$\lambda_0 = \text{vacuum wavelength} = \frac{2\pi c}{\omega}$$

note: θ_d as shown would be a negative value, for this sign convention.

spectral phase difference

$$\Delta \phi(\omega) = \frac{\omega}{c} L_g (n_2 \cos \theta_d - n_1 \cos \theta_i)$$

> where we use the grating eqn to evaluate θ_d

> typically $n_1 = n_2 = \text{air}$

$\theta_i = \text{constant}$ (same for all ω 's)

grating compressor:



reflection geometry.

prism pairs: bases are parallel

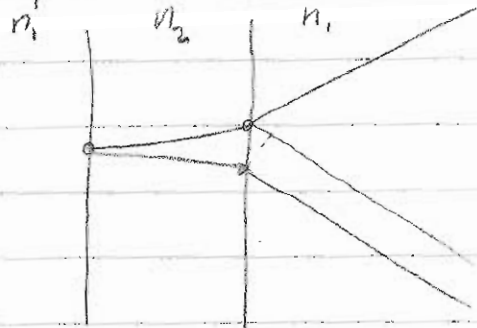


construct as: a tilted window

subtract a "window" of air



Double pass



$$\Delta\phi = 2k_0 n_2 d \cos\theta_r + \pi$$

bc of reflection phase change.

let $\Lambda = 2n_2 d \cos\theta_r$ (optical path)

Fizeau fringes / Newton's rings

r, t : from n_1 to n_2 r', t' : from n_2 to n_1

1st reflection $E_1 = E_0 r$

2nd reflection $E_2 = E_0 t r' t' e^{i k_0 \Lambda}$

$$r' = -r$$

$$t t' = 1 - r^2$$

from Fresnel eqns: see F-P notes.

$$\begin{aligned} I_{out} &\propto |E_1 + E_2|^2 \\ &= I_0 |r|^2 \left| 1 - (1 - r^2) e^{i k_0 \Lambda} \right|^2 \\ &= I_0 |r|^2 (1 + |1 - r^2|^2 - 2(1 - r^2) \cos k_0 \Lambda) \end{aligned}$$

At $k_0 \Lambda = m \cdot 2\pi$ $\cos(k_0 \Lambda) = +1 \rightarrow$ dark fringe.

$k_0 \Lambda = m \cdot 2\pi + \pi \rightarrow$ bright fringe.

IF d is slowly varying $d(x) \rightarrow$ series of fringes.

spacing: $\Lambda = \lambda_0$ total opt. path change = $1 \lambda_0$