

Traveling waves

Note Title

2/8/2008

$$\frac{\partial^2 f}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 f}{\partial t^2} = 0 \quad \text{1D wave eqn.}$$

Let $f = f(x \pm ct)$ or $f(kx \pm \omega t)$

$$\frac{\partial f}{\partial x} = k f'$$

$$\frac{\partial^2 f}{\partial x^2} = k^2 f''$$

$$\frac{\partial f}{\partial t} = \pm \omega f'$$

$$\frac{\partial^2 f}{\partial t^2} = \omega^2 f''$$

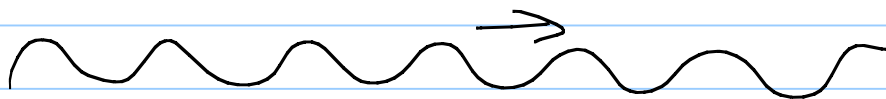
plug into wave eqn.

$$k^2 f'' - \frac{\omega^2}{c^2} f'' = 0 = (k^2 - \frac{\omega^2}{c^2}) f'' = 0$$

true if $k^2 = \frac{\omega^2}{c^2}$, no matter what f .

the shape f moves with "speed" c to the left or right

Eg. $f(x,t) = A \sin(kx - \omega t)$



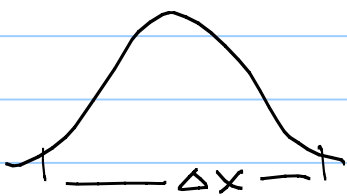
Super position

$$f(x,t) = \sum_i A_i \sin(k_i x - \omega_i t)$$

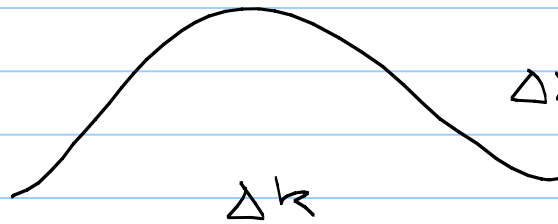
or

$$\int A(k) \sin(kx - \omega t) dk$$

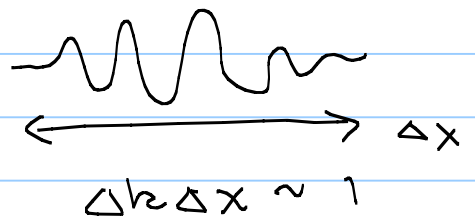
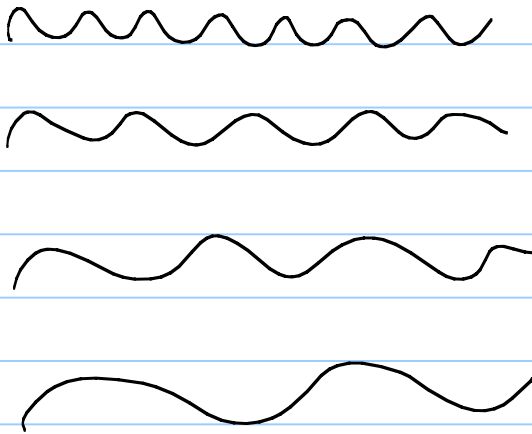
to produce a localized wave packet requires lots of wave numbers



$$\Delta x \Delta k \sim 1$$



$$\Delta x \Delta k \sim 1$$



2 - 8 - 08

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Free particle $V(x) = 0$

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} = E \psi$$

$$\psi'' + \frac{2mE}{\hbar^2} \psi = 0$$

$$\psi'' + k^2 \psi = 0 \quad k^2 = \frac{2mE}{\hbar^2}$$

$$\psi(x) = A e^{ikx} + B e^{-ikx}$$

$$\Psi(x,t) = A e^{i(kx - Et/\hbar)} + B e^{-i(kx + Et/\hbar)}$$

factor out a k $e^{ik(x - \frac{E}{\hbar k} t)}$

$$\frac{E}{\hbar k} = \frac{\hbar k}{2m}$$

$$\Psi(x,t) = A e^{ik(x - \frac{\hbar k}{2m} t)} + B e^{-ik(x - \frac{\hbar k}{2m} t)}$$

plane waves moving to left & right

$$x \pm vt$$

$$v = \frac{\hbar k}{2m}$$

$$k = \pm \frac{\sqrt{2mE}}{\hbar}$$

$$k > 0$$
$$< 0$$

traveling to right
left

de Broglie $p = \hbar k$

$$v = \frac{\hbar k}{2m} = \frac{p}{2m} = \frac{1}{2} v_{\text{classical}}$$

Not normalizable, form a
Linear combination

$$\Psi(x,t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{i(kx - \hbar k^2/2m t)} dk$$

for appropriate choice of ϕ this
can be normalized

now, as usual, we will be
given $\Psi(x, t=0)$ and asked to
find $\Psi(x, t)$.

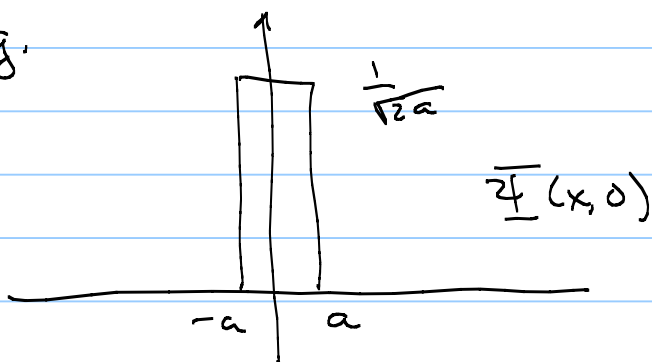
$$\underline{\Psi}(x, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{ikx} dk$$

Hence $\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \underline{\Psi}(x, 0) e^{-ikx} dx$

Recall from 311

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} P(k) e^{ikx} dk \Leftrightarrow P(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{-ikx} dx$$

Eg.



$$\phi(k) = \frac{1}{\sqrt{2\pi}} \frac{1}{\sqrt{2a}} \int_{-a}^a e^{-ikx} dx$$
$$\frac{1}{-ik} e^{-ikx} \Big|_{-a}^a$$

$$\frac{1}{2} \frac{1}{\sqrt{\pi a}} \frac{i}{k} \left[e^{-ika} - e^{+ika} \right]$$

$$-2i \sin ka$$

$$= \frac{1}{\sqrt{\pi a}} \frac{\sin ka}{k}$$

$$\text{So } \psi(x, t) = \frac{i}{\pi} \frac{1}{\sqrt{2a}} \int_{-\infty}^{\infty} \frac{\sin ka}{k} e^{-ik(x - \frac{\hbar k}{2m} t)} dk$$

can't be done analytically

Standard form of a wave packet

$$\psi(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \phi(k) e^{i(kx - \omega t)} dk$$

in our case $\omega = \hbar k^2 / 2m$

$\omega(k)$ is called the dispersion relation

$$\frac{\omega}{k} = v_{\text{phase}}$$

$$\frac{d\omega}{dk} = v_{\text{group}}$$

$$\omega = \frac{\hbar k^2}{2m}$$

for Free particle $\frac{d\omega}{dk} = \frac{\hbar k}{m}$

$$\frac{\omega}{v} = \frac{\hbar k}{2m}$$

$$v_{\text{group}} = v_{\text{classical}} = 2 v_{\text{phase}}$$