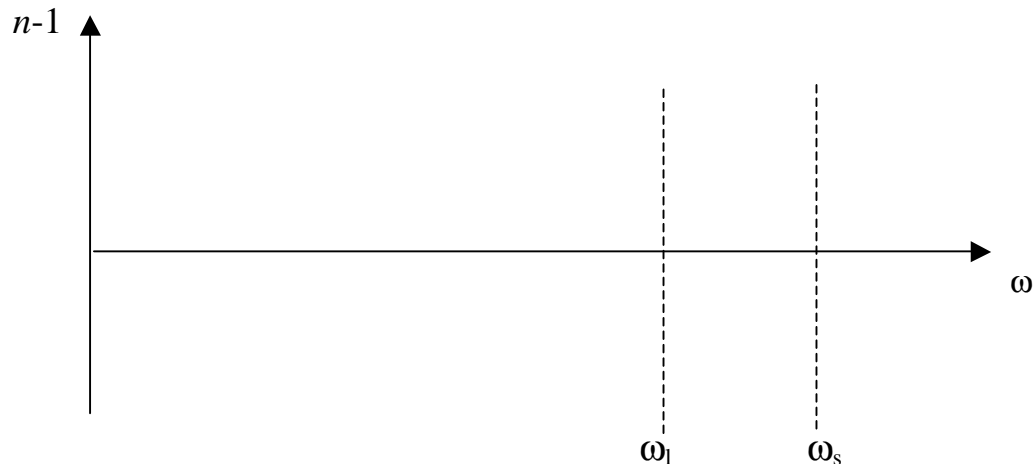


1. Suppose we have a gas of a molecule that has a long axis and **two, equal** short axes. If we apply a static electric field along the x-direction, the molecules will align with the long axis in that direction. We will use a simple, single-resonance model for the response of the molecule to an EM wave. Along the short and long directions of the molecule, the resonance frequencies of the bound electrons are  $\omega_s$  and  $\omega_l$  respectively. Assume the damping coefficient is the same for both resonances, and that the oscillator strengths ( $f_\alpha$ ) are equal to 1. Since the electrons are less tightly bound along the long direction,  $\omega_l < \omega_s$ .
  - a. Consider a plane wave propagating in the z-direction through this gas of molecules aligned in the x-direction. Write two expressions for the complex refractive index ( $n_x$  and  $n_y$ ) experienced by light polarized in the x- and the y-directions. It is not necessary to derive these expressions.
  - b. Sketch two curves for the real part of the refractive index that corresponds to these two resonance frequencies.



- c. Suppose we measure the refractive indices at a wavelength of 500nm and we find that  $n_x - 1 = 5 \times 10^{-4}$  at standard atmospheric pressure, and  $n_y - 1$  is 20% **lower**. The gas is transparent at this wavelength. Estimate values for the two resonance frequencies.
- d. A plane wave of wavelength  $\lambda = 500\text{nm}$ , with linear polarization oriented at  $45^\circ$  to the x-axis is directed into a gas cell that is 10 cm long. **The gas cell has the molecules aligned as described above, and this alignment is unaffected by the input wave.** The pressure is brought up from vacuum to a point where the output polarization is observed to rotate by  $90^\circ$  **compared to the input polarization**. What pressure is required for this to occur? **Treat the gas as ideal and at room temperature.** Be sure to show how you arrive at your answer.

2. Consider a metallic rectangular waveguide filled with a dielectric with an index of refraction  $n$ , and with widths  $a$  and  $b$  along the  $x$  and  $y$  directions, respectively. The propagation direction is in  $z$  (see Fig. 7-5 in Heald and Marion).
- For the TM modes, write an expression for the longitudinal component of the electric field for the guided modes,  $E_z(x, y, z, t)$ . Describe the boundary conditions you are applying for this field component. Define the values of the  $x$ ,  $y$  and  $z$  components of the  $k$ -vector.
  - Show that the lowest-order mode is  $TM_{11}$ .
  - Now assume the dielectric inside the waveguide is a plasma. Calculate an expression for the  $z$ -component for the  $k$ -vector,  $k_z$  expressed in terms of the number density  $N_e$  of the free electrons in the plasma.
  - Calculate an expression for the cutoff frequencies  $\omega_{mn}$  for this plasma-filled waveguide, where  $m, n$  are the mode indices for the  $x$  and  $y$  directions. Show your work and explain the basis for your reasoning.
  - Suppose the waveguide is square, with  $a = b = 0.01$  mm, and the waveguide is empty (vacuum). What is the cutoff wavelength for the lowest mode?
  - If the input wavelength to the square waveguide from part e is  $7\mu\text{m}$ , how many modes does the empty waveguide support? At what number density for the plasma will the waveguide only support a single mode?

3. Consider the radiation damping of a moving, oscillating charge. An electron is launched **at the origin** with a velocity components  $v_{x0}, v_{z0}$  into a region where there is a potential

$$U(x, z) = \frac{1}{2} m \omega_0^2 x^2.$$

- Calculate the initial (**e.g. during the first cycle**) time average radiated power, assuming all velocities are much less than  $c$ , **and that the path of the electron during this cycle is not affected by the energy lost to radiation.**
- Describe how radiation damping affects the motion of the electron: make sketches of the  $x(t)$  and  $z(t)$ . Explain your reasoning.
- Sketch the angular distribution of the radiated power in the low velocity limit  $v \ll c$ , and in the limit where  $v_{z0}/c$  is appreciable. Support your answer by relating this physical situation to other examples in the book.